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# A true triaxial experimental study on porous Vosges sandstone: from strain localization precursors to failure using full-field measurements

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## Abstract

This study systematically investigates the effect of deviatoric loading paths on diffuse and localized deformation developing during the mechanical loading of a high porosity (20%) Vosges sandstone (Eastern France). Laboratory scale experiments are performed using a high pressure true triaxial apparatus, designed to provide access to full-field surface kinematics at high spatial and temporal resolutions during the loading phase. The true triaxial experiments, with independent control of the three principal stress, are conducted at two constant mean stresses, in the brittle-ductile transition regime, and at five prescribed Lode angles, from axisymmetric compression (ASC) to axisymmetric extension (ASE). First, the transition from diffuse towards localized deformation is analyzed in different loading increments and shows an intermediate step of early strain localization, characterized by a large number of early deformation bands developing well before the stress peak and with a predominantly dilatant behavior. Secondly, the evolution of the mechanical behavior and localization patterns, such as deformation band angles and localized dilatancy, indicate a transition from the brittle regime to the ductile regime that is not only dependent on an increase in the mean stress, but also on a decrease in the Lode angle. The analysis of full-field measurements also provides insights into the emergence and evolution of local strains, as deformation structures coalesce or relocate and different failure modes develop depending on the prescribed stress paths.

*Keywords:* true triaxial experiments, strain localization, full-field measurements, localization precursors, Lode angle, sandstone, porous rock

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## 1. General context

The development of strain localization in confined geomaterials is recognized as a fundamental deformation mechanism which generally relates to an evolution in the mechanical response and failure. In high porosity sandstone, this phenomenon is most often apparent in the form of a unique or series of parallel and conjugated deformation bands, characterized by extended structural deformations in narrow and linear zones of concentrated deformation [e.g., 1, 2, 3].

In studies performed under controlled laboratory settings, the manifestation of deformation bands leading to failure in stressed rock samples has been demonstrated to be sensitive to the loading path and the mechanical response of the underlying micro-structure. A well studied phenomenon in high porosity rocks is the transition into a ductile regime with an increase in the first invariant of the stress tensor, i.e. the mean stress [e.g., 4, 5]. The conventional triaxial apparatus, where a stress deviator is applied in the major principal direction while the confining pressure is applied on a jacketed cylindrical specimen, is well suited to assess the role of the mean stress. In such experimental conditions, the apparatus-sample contact boundaries are well defined, by the superposition of a confining fluid and a deviatoric loading using an axial unidirectional rigid piston. Nonetheless, the lack of independent control between the radial stresses limits the imposed stress paths to axisymmetric compression and extension conditions, leaving an intermediate broad spectrum of the octahedral (deviatoric) plane unexplored.

To circumvent this limitation, different variations of the true triaxial apparatus (TTA) have been developed in various experimental laboratories, aiming to investigate the role of alternative stress paths to the classical axisymmetric cases [6]. The most prominent method for mechanical testing of rocks in true triaxial conditions consists in performing loading experiments at different minor and intermediate principal stress levels, while the major principal stress evolves with incremental axial strain. As such, the design advancements of the TTA have enabled to demonstrate the important role of the intermediate principal stress on the deformation and failure mechanisms in various types of rocks. In terms of the resulting stress invariants during this type of loading, while intermediate values of the Lode angle, between  $0^\circ$  and  $60^\circ$ , can be achieved, the Lode angle continuously and systematically decreases in conjunction with an increase in the mean stress ( $\sigma_m$ ). Therefore, the independent effect of the stress invariants on the evolution of the

38 material response cannot be directly assessed. Additionally, their continu-  
39 ous and simultaneous evolution during loading, inherent to principal stress  
40 controlled loading, results in the distribution of stress states at failure to be  
41 scattered in certain regions of the meridian and octahedral planes. Figures  
42 1a.-l. illustrates this effect for selected data sets retrieved from the literature  
43 for principal stress controlled testing of various rock samples, where values  
44 of the peak octahedral stress, as defined in equation 1, are represented as  
45 projections on a single octahedral plane.

46 An alternative experimental approach consists in systematically investi-  
47 gating different stress paths with prescribed invariants of the stress tensor,  
48 by imposing an incremental relation between the principal stresses. Ex-  
49 periments performed under this type of loading procedure provide valuable  
50 insight on the independent effect of isotropic and deviatoric loading paths on  
51 the mechanical response and deformation mechanisms. Ingraham et al. [13]  
52 performed a series of experiments under such loading conditions on samples  
53 of a high porosity sandstone, by imposing two independent invariants during  
54 the loading phase: the mean stress ( $\sigma_m$ ) and the Lode angle ( $\theta$ ), as defined  
55 in equation 1. The peak stress recorded for each test are thus organized on  
56 selected meridian and deviatoric planes (see Figure 1.m). Ma et al. [14] have  
57 followed a similar methodology, for a mixed loading mode at a constant minor  
58 principal stress and constant Lode angle in two different porous sandstone  
59 samples. The results from these two studies have consistently shown a mea-  
60 surable effect of the Lode angle on failure and localization patterns analyzed  
61 in post-mortem observations. Nonetheless, experimental data on the subject  
62 are scarce and many questions remain open as to the mechanisms involved  
63 during the deformation process of porous rocks subjected to general stress  
64 paths.

65 It is thus essential to further study the effect of different loading condi-  
66 tions in terms of the mechanical response, but also in their relation to the  
67 emergence and development of localization structures in porous rocks. The  
68 purpose of the present experimental campaign is therefore to systematically  
69 explore the effect of different invariant controlled stress paths on the me-  
70chanical response, combined with a unique full-field measurement technique  
71 to characterize the diffuse and localized deformation in laboratory samples of  
72 a Vosges sandstone, a high porosity quasi-isotropic model sedimentary rock.

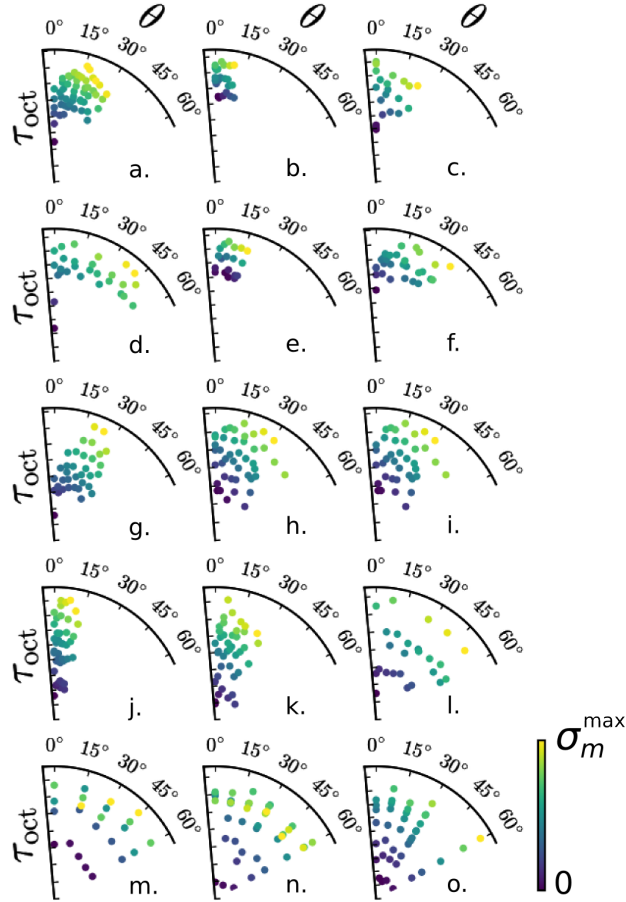


Figure 1: Representation of stress state at failure in the octahedral plane, arranged by experimental campaigns retrieved from the literature: Dunham dolomite (a), Inada granite (b), Manazuru andesite (c), Mizuho trachyte (d), Orikabe monzonite (e), Solhofen limestone (f), Yamaguchi marble (g) [from 7]; Yuubari shale (h), Shirahama (i) [8] [data from 9]; Westerly granite (j) [10]; KTB amphibolite (k) [11]; Taiwan siltstone (l) [12]; Castle-gate sandstone (m) [13]; Coconino sandstone (n), Bentheim sandstone (o) [14]. The color scale indicates the mean stress level reached at the deviatoric stress peak, the value of the maximum mean stress being specific to each study.

## 73 2. Methods and material

### 74 2.1. Experimental setup

75 The high pressure true triaxial apparatus (TTA) used for this experimen-  
76 tal campaign has been designed and assembled at Laboratoire 3SR (Grenoble)  
77 for the specific purpose of testing rock samples in biaxial and true tri-  
78 axial loading conditions [15]. The apparatus is designed to accommodate  
79 prismatic samples of  $50 \times 30 \times 25 \text{ mm}^3$  inside a specially fabricated soft mem-  
80 brane installed in the main confining chamber. The loading mechanism of  
81 this apparatus relies on a combination of direct fluid pressure and a pair of  
82 rigid pistons to apply stresses at the surfaces of the prismatic samples. The  
83 mixed boundary conditions aligned with the surfaces of the sample are thus  
84 imposed by a combination of rigid stainless steel and sapphire glass contacts  
85 in the major (1) and intermediate (2) directions, respectively; and a soft  
86 contact through the jacket in the minor (3) direction. The loading device  
87 is rated for a maximum isotropic confinement of up to 100 MPa, applied by  
88 increasing the fluid pressure inside the confining chamber hosting the iso-  
89 lated sample. The stress deviators, decoupled in the major and intermediate  
90 directions, are applied by the means of pressure controlled hydraulic pistons  
91 (self-equilibrated with respect to the cell pressure). In the major direction  
92 1, a set of sliding pedestals transfer the load on the surface of the specimen  
93 while accommodating displacements in the orthogonal direction. For the  
94 specified sample dimensions and through the effect of the balanced ram, the  
95 stress deviator, in both the major and intermediate directions, can reach up  
96 to 670 MPa and 530 MPa, respectively. A detailed description of the ap-  
97 paratus is given in [16]. A recent modification consist in the manufacturing  
98 of urethane molding rubber to produce soft reusable jackets, making them  
99 adaptable to strain gauge installation by passing cable through sealed con-  
100 duits; an essential improvement to obtain direct surface strain measurements  
101 in the out-of-plane (non-visible) direction 2 [17].

### 102 2.2. Loading Procedure

103 Using linear relationships between the principal stress values ( $\sigma_1$ ,  $\sigma_2$  and  
104  $\sigma_3$ ), a loading procedure was developed for the control system of the appa-  
105 ratus to perform true triaxial loading experiments with imposed loading paths  
106 based on the invariants of the Cauchy stress tensor. While any combination  
107 of mutually independent set of invariants are suitable to objectively define  
108 general stress state increments, the Octahedral-Lode invariants is preferred

109 as it describes a convenient polar coordinate system around the trisectrix, or  
 110 isotropic stress state line, in the principal stress space. The selected invari-  
 111 ants are thus defined as

$$\begin{aligned}
 \sigma_{av} &= \frac{1}{3}[\sigma_1 + \sigma_2 + \sigma_3] \\
 \tau_{oct} &= \frac{1}{3}[(\sigma_1 - \sigma_2)^2 + (\sigma_1 - \sigma_3)^2 + (\sigma_2 - \sigma_3)^2]^{1/2} \\
 \theta_\sigma &= \arctan \left[ \sqrt{3} \frac{\sigma_2 - \sigma_3}{(\sigma_1 - \sigma_2) + (\sigma_1 - \sigma_3)} \right]
 \end{aligned} \tag{1}$$

112 where  $\sigma_{av}$  is the average stress, and  $\tau_{oct}$  and  $\theta$  are respectively the octahedral  
 113 stress and Lode angle, measuring independently the deviatoric stress mag-  
 114 nitude and orientation in the octahedral plane (see Figure 2). The strain  
 115 invariants ( $\epsilon_{av}$ ,  $\gamma_{oct}$ ,  $\theta_\epsilon$ ) are equivalently defined in terms of the principal  
 116 values of the strain tensor. In the adopted solid mechanics convention, com-  
 117 patible with the presented full-field measurements, both principal stresses  
 118 and strains are positive in tension. To simplify the discussion of results, the  
 119 invariants along the trisectrix of the principal stress and strain space are pre-  
 120 sented in terms of the mean stress,  $\sigma_m = -\sigma_{av}$ , and the volumetric strain,  
 121  $\epsilon_v = 3\epsilon_{av}$ . Therefore  $\sigma_m$  is positive in compression and  $\epsilon_v$  is positive during  
 122 dilatancy. The Lode angle, as defined in equation 1, spans from  $0^\circ$  in axisym-  
 123 metric compression (ASC,  $\sigma_2 = \sigma_3$ ) to  $60^\circ$  in axisymmetric extension (ASE,  
 124  $\sigma_1 = \sigma_2$ ). It is a measure of the deviatoric stress orientation in the stress  
 125 space, and can be intuitively interpreted as an indication of the intermediate  
 126 principal stress magnitude, with respect to the magnitude of the minor and  
 127 major principal stresses at a given mean stress.

128 For the presented experimental campaign, a consistent procedure was  
 129 followed for each experiment performed at prescribed loading paths: (1) a  
 130 sample alignment phase is initiated to insure centering of the contact between  
 131 the pistons, pedestals and the sample surfaces in directions 1 and 2; (2) a  
 132 subsequent isotropic monotonic loading phase consists of increasing the cell  
 133 pressure at a rate of 2 MPa/min, until the prescribed mean stress is reached;  
 134 (3) finally, a deviatoric loading phase consists in monotonically increasing the  
 135 octahedral stress while maintaining  $\sigma_m$  and  $\theta$  at constant values. During this  
 136 last loading phase, a PID controller feedback loop prescribes the hydraulic  
 137 pressure to update the intermediate and minor principal stresses, in terms of  
 138 the major principal stress, according to the relationship:

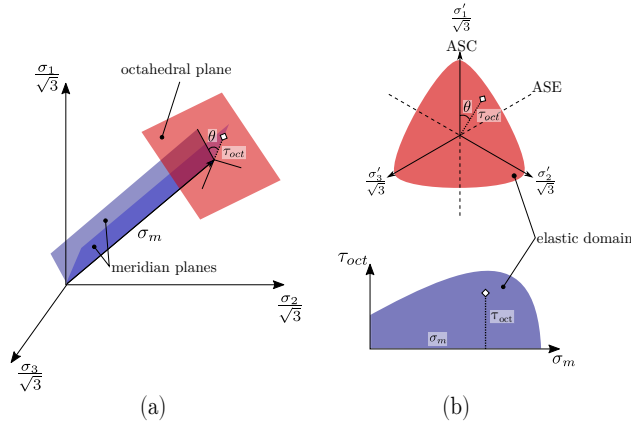


Figure 2: A graphical representation of the cylindrical Lode coordinate system in the principal stress space (a), based on the three invariants  $\sigma_m$ ,  $\tau_{oct}$  and  $\theta$ . In (b), the octahedral plane (normal projection on a plane perpendicular to the isotropic state line) and the meridian plane ( $\tau_{oct}$  vs.  $\sigma_m$ ) provide complementary 2D information of the stress state.

$$\begin{aligned}
 \sigma_2 &= \frac{3 \tan \theta - \sqrt{3}}{2\sqrt{3}} (\sigma_1 - \sigma_m) + \sigma_m \\
 \sigma_3 &= \frac{-3 \tan \theta - \sqrt{3}}{2\sqrt{3}} (\sigma_1 - \sigma_m) + \sigma_m
 \end{aligned} \tag{2}$$

139 obtained from a simple algebraic rearrangement of equations 1. The hy-  
 140 draulic syringe pumps are controlled based on this set of equations, while  
 141 including additional proportionality parameters to calibrate the hydraulic  
 142 pressure in the piston according to the sample surface stresses and frictional  
 143 correction terms. These linear correction terms take into account the varia-  
 144 tion in the dynamic friction effect between the piston shaft and the seals with  
 145 the variation in hydraulic pressure in the main confining chamber, which is  
 146 equal to  $\sigma_3$ . The correction of  $\sigma_1$  ranges from 4.6 MPa to 9.7 MPa and the  
 147 correction  $\sigma_2$  from 3.2 MPa to 9.3 MPa, for values of  $\sigma_3$  between 0 MPa and  
 148 -90 MPa.

149 Due to significant non-linear elastic deformation of the apparatus during  
 150 the deviatoric loading phase, the vertical strain rate could not be explicitly  
 151 prescribed based on the measurements from the external LVDT linked to  
 152 the control system. The exact strain measurements are therefore a priori



Table 1: Mineral characterization of Vosges sandstone

mineral composition	93% quartz 5% microcline 2% mica and kaolinite
porosity	21%
grain size	150-450 $\mu\text{m}$
mean grain size	300 $\mu\text{m}$

153 unknown during the experiments, since applied corrections from the sample  
 154 boundary displacement are obtained by digital image correlation during the  
 155 post-experimental analysis. Instead, the monotonic loading in the major  
 156 principal direction is controlled at a constant injection rate (1.5 cc/hr) in  
 157 the vertical piston chamber, resulting in a mean vertical strain rate for all  
 158 the experiments of  $4.2 \times 10^{-7} \text{ s}^{-1}$  at the beginning of the deviatoric loading  
 159 phase and  $2.3 \times 10^{-6} \text{ s}^{-1}$  at the end of the hardening phase. This strain rate  
 160 evolution during the experiment is a function of the compressibility of the  
 161 injection fluid, an hydraulic oil with entrapped air pockets, and the stiffness  
 162 of the true triaxial apparatus.

### 163 *2.3. Material*

164 A total of ten samples has been prepared using a single block of Vos-  
 165 ges sandstone originally retrieved from the Woustwiller quarry, in Eastern  
 166 France. The mineral composition and micro-structural characterization of  
 167 this sedimentary rock, reported here in Table 1, were documented by Bésuelle  
 168 et al. [18] in an earlier experimental study on this particular porous sand-  
 169 stone. The initial porosity for each of the retrieved sample has been measured  
 170 at  $21\% \pm 1$ , showing a satisfactory consistency between sampling locations  
 171 of the block.

172 The prismatic samples, of approximate size  $50 \times 30 \times 25 \text{ mm}^3$ , are pre-  
 173 pared as to have symmetric boundary conditions on opposite surfaces of the  
 174 geometry by performing successive iterations of machining, polishing and  
 175 later applying watercolor paint coatings. The allowable tolerance on the  
 176 parallelism is  $\pm 0.02 \text{ mm}$  for each pair of surfaces in contact with the hard  
 177 platens (hardened steel in direction 1 and sapphire glass in direction 2) and  
 178  $\pm 0.03 \text{ mm}$  for the pair of surfaces in contact with the urethane membrane, af-  
 179 ter the strain gauge installation. The surfaces in contact with the platens are  
 180 finally lubricated using a thin coating of a stearic acid and vaseline mixture,

181 effectively minimizing end friction effects under high confining pressures [19].  
 182 This thin lubricant coating being translucent, the visual aspect of the fine  
 183 speckle pattern applied over the surface is maintained after an homogeneous  
 184 contact is established between the sample and the sapphire glass.

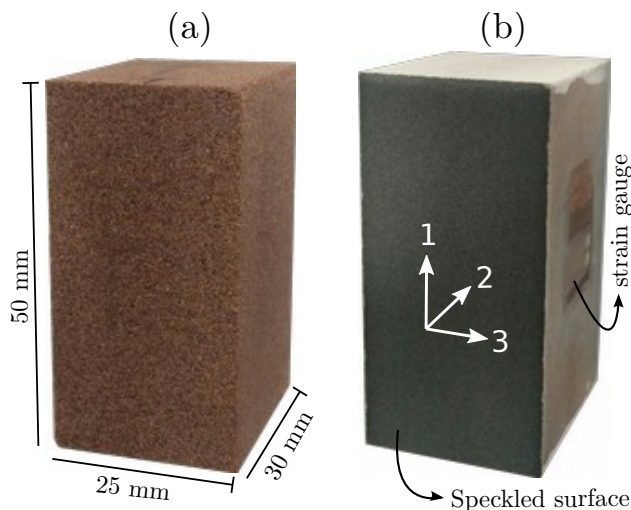


Figure 3: Intact prismatic samples of Vosges sandstone showing (a) the natural porous surfaces and (b) prepared surfaces with horizontal strain gauges, painted surfaces and the monochromatic speckle. Orthogonal vectors mark the spatial directions associated with the principal stress space.

#### 185 2.4. Instrumentation and measurements

186 The unique conception of the TTA used in this study enables the moni-  
 187 toring of full-field surface kinematics, using digital image correlation (DIC),  
 188 from photographs of one of the sample surface during the loading phase. The  
 189 out-of-plane strains are also measured directly on the sample surface using  
 190 strain gauge instrumentation.

191 The full-field photographs of the sample surface are captured through  
 192 a sapphire viewport, oriented in the 1-3 plane (orientation as in Figure 3).  
 193 The digital imaging system, a single-lens reflex (DSLR) 36 Mpx Nikon D800E  
 194 camera, provides a spatial pixel size of 7 microns over the entire surface of  
 195 the sample. In the camera configuration used to minimize noise and shutter  
 196 vibrations, photographs were taken at regular intervals with a minimum time  
 197 period of 3 seconds/image.

198 On the white painted and polished visible surface of the sample, a fine  
199 monochromatic speckle pattern is applied by paint atomization, using a pre-  
200 cision airbrush, to create a random deposition of black watercolor droplets  
201 over the whole region of interest (ROI) (Figure 3b). This process enhances  
202 the optical texture and uniqueness of local zones of interest (ZOI) from the  
203 natural texture created by the surface heterogeneity on the surface of the  
204 porous sandstone. The DIC algorithm relies on the conservation of optical  
205 flow during the deformation process. Therefore, the spatial resolution, con-  
206 vergence and systematic error in the process rely on the quality of the optical  
207 texture in each ZOI. The texture quality is assessed for each created speckles  
208 by computing the autocorrelation function [20]. The autocorrelation radius  
209 is a scalar measurement of the pixel displacement where the autocorrelation  
210 function of a ZOI reduces by half, indicating the level of self similarity of  
211 the texture around this area [21, 22]. Measurements of the autocorrelation  
212 radius in the vertical and horizontal directions are shown in Figure 4a and  
213 4b respectively for a typical surface ROI. For DIC, a speckle pattern with an  
214 autocorrelation radius of around 2 pixels is considered to provide an optimal  
215 texture quality which is independent of the acquisition noise.

216 The DIC algorithm used to performed the image analysis is the open  
217 source Software for Practical Analysis of Materials (Spam) developed as a  
218 collaborative and open source project [23]. The correlation algorithm, de-  
219 signed for 2D and 3D analysis, is based on a two step approach: an initial  
220 correlation to the nearest pixel, generating an approximation of the defor-  
221 mation field by tracking large rigid body motions, and subsequent iterations  
222 consisting in a sub-pixel correlation based on the initial approximation ob-  
223 tained from the first step. This latter operation allows for all possible affine  
224 (six degrees of freedom in 2D) transformations of the ZOI, to accurately track  
225 the local displacement at each node. One of the purposes of the displacement  
226 field obtained from DIC is to apply corrections to the sample’s boundary  
227 displacement in the 1-3 visible plane, providing accurate global strain mea-  
228 surements. The strain in the horizontal direction 3 is computed from the  
229 horizontal displacement of DIC points at the two vertical boundaries. In the  
230 vertical direction 1, the external LVDT measurement is corrected at regular  
231 intervals using accurate DIC measurements, compensating deformation ar-  
232 tifacts from the compliance of different parts of the apparatus (loading cell,  
233 surface contact interfaces) measured externally. The contribution of in-plane  
234 deformation to the global volumetric deformation of the sample is calculated  
235 from the sum of averaged outward displacements at each boundary. This

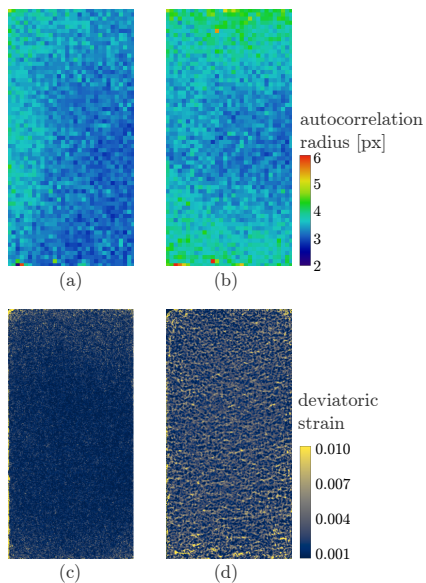


Figure 4: Assessment of speckle quality and measurement error in sample G16, at  $\sigma_m = 60$  MPa and  $\theta = 0^\circ$ , showing: the fields of autocorrelation radii in the vertical (a) and horizontal (b) direction, and image correlation of two photographs in the same deformation state (c), compared to the increment 1-2 reported in Figure 7.

236 method is equivalent to the computation of the contour integral for small  
 237 deformation gradients.

238 Additionally, the two-dimensional fields of incremental local strain, cal-  
 239 culated from the spatial first order derivative of the displacement field, are of  
 240 clear interest in identifying localization patterns and their evolution. Nonethe-  
 241 less, due to reduced continuity compared to the discrete displacement field,  
 242 the local strain field is more sensitive to the errors in the image acquisition  
 243 process and selected DIC parameters (acquisition noise, choice of ZOI size,  
 244 convergence error resulting from the DIC process). It is thus necessary to ap-  
 245 ply local corrections and filtering of the displacement field to obtain a better  
 246 representation of the derived strain field. First, the field of correlation resid-  
 247 uals, quantifying the correlation error and non-linear deformation, is used  
 248 to mask the displacement field; the values at poorly correlated nodes are  
 249 replaced by a weighted interpolation of surrounding converged nodes. Sec-  
 250 ond, a sequential median and mean convolution of the nearest neighbors  
 251 is performed to render a continuous field by eliminating isolated singularities,  
 252 at the expense of degrading the local pixel-scale deformation measurements.

253 From the resulting filtered displacement field, a local strain tensor is cal-  
254 culated based on the displacement of 4 neighboring nodes using a regular  
255 grid layout and linear finite element shape function interpolations. This pro-  
256 cedure yields a spatial map of the local deformation gradient tensor at the  
257 center of each cell. Since large rotations and strains may develop locally, large  
258 transformation formalism is used for the local deformations derived from the  
259 incremental displacement field (see Appendix A).

260 To assess the effect of acquisition and correlation errors on the strain  
261 measurements, Figure 4c shows a deviatoric strain field calculated using two  
262 photographs acquired in the same deformation state, at the beginning of  
263 the deviatoric loading phase (at photograph 1). In the deformation range of  
264 interest, the strain measurement error is seen to be only slightly more im-  
265 portant in regions of the sample surface where the autocorrelation radius  
266 is higher. However, this error remains negligible in comparison to the mag-  
267 nitude of localized deformation observed in selected subsequent increments  
268 of the deviatoric loading phase (increment 1-2 is illustrated in Figure 4d for  
269 comparison).

270 To measure the out-of-plane deformation in direction 2, strain gauges were  
271 installed horizontally at mid-height on both sides of the sample surfaces in  
272 contact with the soft membrane. For the axisymmetric compression experi-  
273 ments, where  $\sigma_2 = \sigma_3$ , it was observed that strain gauges in direction 2 and  
274 DIC measurements in direction 3 were in good agreement until pronounced  
275 localization effects start to develop. This reveals a good correspondence be-  
276 tween different measuring techniques as well as strain homogeneity despite  
277 different types of boundary conditions with different end-friction and possi-  
278 ble bulging effects developing in independent directions 2 and 3. The global  
279 volumetric strain is represented as to take into account the full three dimen-  
280 sional in-plane and out-of-plane components in each of the three principal  
281 strain directions.

### 282 3. Experimental Results

283 Guided by the study of Bésuelle et al. [18] in which axisymmetric com-  
284 pression (ASC) and extension (ASE) loading experiments were performed  
285 on samples of the Vosges sandstone (extracted from the same initial block),  
286 the range of mean stress was selected to investigate deformation mechanisms  
287 associated with the brittle-ductile transition regime. The present experimen-  
288 tal campaign thus consisted of ten monotonic loading experiments following

Table 2: Summary of stress states and angle  $\beta$  at  $\tau_{oct}^{peak[1]}$  for each of the 10 loading paths

$\sigma_m^{peak}$ [MPa]	$\theta^{peak}$ [°]	$-\sigma_1^{peak}$ [MPa]	$-\sigma_2^{peak}$ [MPa]	$-\sigma_3^{peak}$ [MPa]	$\tau_{oct}^{peak}$ [MPa]	$\beta$ [°]
61.1	0.0	128.7	27.2	27.3	47.8	50
61.7	16.0	121.0	46.8	17.3	43.6	56
62.3	29.6	110.6	61.9	14.5	39.3	61
61.9	45.4	96.9	75.1	13.7	35.2	64
62.0	59.9	87.9	87.8	10.4	36.5	64
90.0	0.0	169.0	50.4	50.4	55.9	43
90.1	15.1	166.4	69.8	34.0	55.9	47
90.0 <sup>[2]</sup>	14.9 <sup>[2]</sup>	168.7	68.9	32.6	57.5	-
91.4	32.4	148.8	94.2	31.2	48.1	54
90.0	44.9	136.9	107.1	26.2	46.8	56
89.9	59.9	122.8	122.7	24.1	46.5	56

[1] The *peak* superscript designates the moment when the octahedral stress reaches its maximum value.

[2] Duplicated experiment at  $\sigma_m=90$  MPa and  $\theta=15^\circ$  not represented in subsequent Figures.

289 invariant controlled loading paths at five prescribed Lode angles (0, 15, 30,  
290 45 and 60°) at two mean stresses (60 and 90 MPa).

291 A summary of the peak principal stresses and corresponding stress invari-  
292 ants is presented in Table 2. In this Table, a duplicated data point at  $\sigma_m=90$   
293 MPa and  $\theta=15^\circ$  is provided. Since only the stress values for this additional  
294 experiment are available, it is not represented in subsequent Figures, but  
295 helps to consolidate the result for this loading path. Due to differences in  
296 the correction terms applied to Equation 2, embedded in the servo-controlled  
297 loading program (constant for tests at  $\sigma_m = 60$  MPa; linear dependency with  
298  $\sigma_3$  for tests at  $\sigma_m = 90$  MPa) and in the post-experimental correction func-  
299 tions (linear dependency with  $\sigma_3$  for all tests), there is a slight discrepancy  
300 between the mean stress and Lode angle initially prescribed and the peak  
301 values of the controlled invariants reported in Table 2. The importance of  
302 the correction terms dependence on the cell pressure ( $\sigma_3$ ) was only assessed  
303 after the first part of the experimental campaign. The deviation is however  
304 not significant compared to the mean stress and Lode angle intervals between  
305 the different selected stress paths. Therefore, the closest prescribed values  
306 will be used to reference the different loading paths.

307 *3.1. Mechanical response*

308 The evolution of global octahedral stress ( $\tau_{oct}$ ) and volumetric strain ( $\epsilon_{vol}$ )  
309 during the deviatoric loading phase is reported for each experiment against  
310 the octahedral strain ( $\gamma_{oct}$ ) in Figure 5. A first comparison of the different  
311 mechanical responses at the two studied mean stresses indicates that, irre-  
312 spective of the Lode angle prescribed for each experiment, an increase in the  
313 mean stress, from 60 MPa to 90 MPa, favours a slower bending of the octa-  
314 hedral stress versus strain curve during the pre-peak regime, while the initial  
315 stiffness is essentially the same. Consequently, the peak octahedral stress is  
316 found to increase with the mean stress, a classical result often reported in the  
317 literature for experiments following ASC loading paths. While the increase  
318 in the peak stress with the mean stress is more pronounced at a Lode angle  
319 of  $15^\circ$ , the variation in this increase with the Lode angle remains moderate  
320 and appears to become relatively constant between experiments at higher  
321 Lode angles of  $30^\circ$  to  $60^\circ$  (Figure 6).

322 At each of the two investigated mean stresses, loading paths following  
323 different Lode angles in the octahedral plane also promote a significant vari-  
324 ation in the mechanical response during the late pre-peak regime. During  
325 this stage of loading, the octahedral stress level along the stress-strain curves  
326 is consistently lower at higher Lode angles. A noticeable difference in the  
327 stress-strain response is also apparent between the lower ( $0^\circ$ ,  $15^\circ$ ) and higher  
328 ( $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ ) values of the Lode angle (mostly apparent at  $\sigma_m=90$  MPa).  
329 This marked difference in the octahedral stress evolution during the pre-peak  
330 regime results in a consistent decrease in the peak octahedral stress with in-  
331 creasing Lode angle, while the octahedral strain corresponding to the peak  
332 stress tends to increase. Between the two prescribed mean stresses, this in-  
333 fluence of the Lode angle on the peak stress occurs over a different range of  
334 Lode angles. In Figure 6a, it is seen to steadily decrease from  $0$  to  $45^\circ$  at  
335  $\sigma_m=60$  MPa and over a shorter range of  $15^\circ$  to  $30^\circ$  at  $\sigma_m=90$  MPa (con-  
336 solidated by the additional data point in Table 2). The non-linearity in the  
337 peak octahedral stress transition indicates an evolution of the shape of the  
338 failure surface with both the mean stress and the Lode angle (Figure 6b).  
339 Furthermore, the variation in the peak octahedral stress over the full range  
340 of Lode angles, from  $0^\circ$  to  $60^\circ$ , is of comparable magnitude to the effect of  
341 the mean stress for the range investigated in the present study (from 60 to  
342 90 MPa). As for the residual stress, after peaking and softening, the values  
343 depend on the mean stress and are rather independent of the Lode angle,  
344 except for the tests at  $0^\circ$  angle where the residual stress is clearly above the

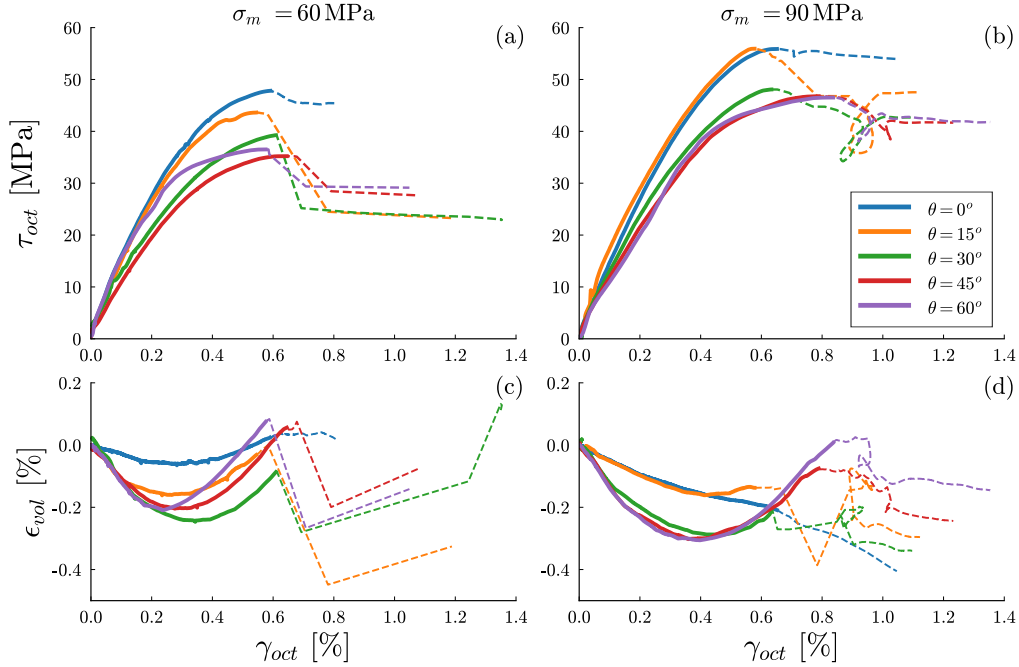


Figure 5: Octahedral stress ( $\tau_{oct}$ ) and volumetric strain ( $\epsilon_{vol}$ ) as a function of octahedral shear strain ( $\gamma_{oct}$ ), analogue to  $\tau_{oct}$  in the strain space. The curves are dashed in the post peak regime, indicating a loss of quasi-static conditions and apparatus displacement control regulation in all but tests at a Lode angle of  $0^\circ$ .

345 other values.

346 The evolution of global volumetric deformation, combining measurements  
 347 using DIC in the 1-3 plane and strain gauges in the out-of-plane direction  
 348 2, is also affected by the prescribed mean stress and Lode angle. In this  
 349 particular case, for selected loading paths at constant mean stresses, it can  
 350 be noted that the elastic part of the volumetric deformation is theoretically  
 351 vanishing. Therefore, the initial compactant behavior, represented in Figure  
 352 5c-d, suggests an inelastic deformation is occurring at the initiation of the devi-  
 353 ator loading phase. At the mean stress of 60 MPa, all samples displaying  
 354 an initially compactant behavior transition into a dilatant phase at approx-  
 355 imately halfway through the pre-peak phase. At this lower mean stress, a  
 356 distinctly lower volumetric change is observed for the Lode angle of  $0^\circ$ . In  
 357 comparison, the samples tested at higher Lode angles display both a more  
 358 pronounced initial compactant and a subsequent well defined dilatant phase.



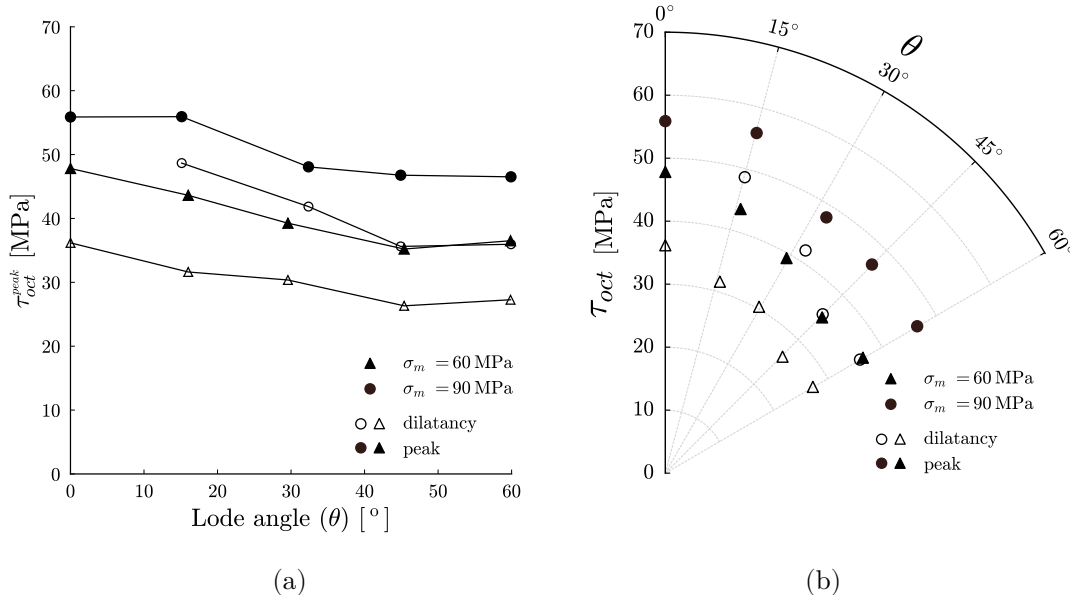


Figure 6: Evolution of the octahedral stress at dilatancy ( $\tau_{oct}^{dilat}$ ) and peak stress ( $\tau_{oct}^{peak}$ ) in (a). The same dataset is represented in the octahedral plane in (b).

359 At the mean stress of 90 MPa, the volumetric deformation, for the different  
 360 loading paths, follows two distinct evolutions during the pre-peak regime: at  
 361 the lower Lode angles of  $0^\circ$  and  $15^\circ$ , a monotonic volumetric deformation  
 362 sustains a moderately compactant behavior of the samples (however, a very  
 363 slight dilatation is observed at  $15^\circ$  towards the stress peak); whereas, at the  
 364 higher Lode angles, the response is similar to the behavior observed at the  
 365 lower mean stress, displaying a more pronounced compactancy followed by a  
 366 well defined dilatant phase. It is possible to identify the octahedral stress at  
 367 the octahedral strain where the inception of dilatancy occurs for each loading  
 368 path. In Figure 6, the decrease in the octahedral stress at dilatancy with de-  
 369 creasing mean stress and increasing Lode angle is concordant with the trend  
 370 in the transition of the peak octahedral stress. In this Figure, the value at  
 371  $\sigma_m = 90$  MPa and  $\theta = 0^\circ$  is omitted since the volumetric deformation is con-  
 372 tinuously compactant for this loading path. At the higher mean stress and  
 373 intermediate values of the Lode angle, the difference between the octahedral  
 374 stress at dilatancy and at the peak decreases with decreasing Lode angle,  
 375 while it remains relatively constant for all Lode angles at the lower mean  
 376 stress.

377 In the post-peak regime, all samples tested at a mean stress of 90 MPa  
378 have a compactant total volume strain past their respective peak octahedral  
379 stress. However, since mature localization structures tend to emerge around  
380 the peak stress, global measurements from average boundary displacements  
381 become discontinuous in this phase of the mechanical response. The global  
382 strain measurements are therefore less representative of the predominant de-  
383 formation mechanism occurring inside the sample. In addition to being the  
384 dominant deformation mechanism at the sample scale during the post-peak  
385 regime, localized deformations also play an important role on the mechan-  
386 ical response during the pre-peak regime. In the next section, incremental  
387 full-field measurements are presented and compared for the different loading  
388 paths, enabling to characterize the emergence of precursors and early de-  
389 formation structures during the pre-peak regime, and their relation to the  
390 development of highly localized structures near and past the peak stress.

### 391 *3.2. Localization patterns*

392 A full-field representation of local deformation, obtained from DIC in  
393 the 1-3, plane enables the identification of strain localization structures de-  
394 veloping during the deviatoric loading phase. In the multi-page Figure 7,  
395 incremental 2D deviatoric and volumetric strain fields for each of the 10 ex-  
396 periments are depicted in consecutive series, from photograph  $n \rightarrow n + 1$ .  
397 Each increment in the series, associated with each experiment at different  
398 loading paths, represents different characteristic phases of the mechanical  
399 and kinematic behavior: an early quasi-linear phase (1-2), an intermediate  
400 phase of volumetric compaction (2-3), the inception of dilatancy generally  
401 associated with the inception of early deformation bands (3-4), the intensifi-  
402 cation of localized zones into well developed early deformation bands (4-5),  
403 the development of mature localization structures organized into a single or  
404 multiple parallel and conjugated bands (5-6), brittle faulting or propagation  
405 of mature deformation band(s) (6-7). The series of photographs, from which  
406 the represented incremental strain fields are produced, are labeled on the  
407 associated individual stress-strain and volumetric curves.

408 In the initial volumetric compaction phase represented in increment 1-2,  
409 the deviatoric strain fields for all loading paths show relatively similar pat-  
410 terns of weakly defined regions of higher strain deformation. During this early  
411 loading phase, the deviatoric strain fields thus present some identifiable het-  
412 erogeneities at the finer scale, where a spatial variation in the deformation is  
413 noticeable between neighboring correlation points forming small distributed

414 clusters. Nonetheless, the represented strain fields do not, at this stage, ex-  
415 hibit any strong evidence of localization clustering or organization at the  
416 sample scale.

417 In full-field images of subsequent increments, following photograph 2 in  
418 respective experiments, two distinct types of deformation bands can be iden-  
419 tified to sequentially emerge into well defined localization structures. The  
420 first type, referred to as *early deformation bands*, is observed to emerge well  
421 before the peak stress and develop in series of parallel and conjugated linear  
422 clusters of limited length, visible mostly in the central section of the ROI.  
423 The localization structures which are well defined in increments 3-4 and 4-5,  
424 can also be observed, if one looks closely, as weak precursors in increment  
425 2-3. The second type, referred to as *mature deformation bands*, later develop  
426 close to the peak stress and is characterized by one or multiple extended lo-  
427 calized regions traversing across the sample. Mature deformation bands can  
428 be seen to initiate during increment 4-5 and 5-6, and to fully develop during  
429 increments 5-6 and 6-7 in respective experiments.

430 The early deformation bands are seen to emerge from a progressive clus-  
431 tering of strain localization in specific areas of the visible surface of the  
432 samples, forming preferentially orientated linear structures which are contin-  
433 uously evolving throughout the later pre-peak regime. During this phase, the  
434 deviatoric strain inside early deformation bands gradually intensifies and fur-  
435 ther localizes in specific regions of the ROI. This evolutive behavior forming  
436 localized regions is most noticeable for experiments following loading paths  
437 at the higher mean stress of 90 MPa and higher Lode angles. It is also dur-  
438 ing increments 3-4 and 4-5, following the inception of the incremental global  
439 dilatancy, that the volumetric strain fields start to show significant localized  
440 dilatant zones, preferentially organized where localization structures are also,  
441 and more noticeably, observed in the deviatoric strain fields. It is therefore  
442 worth noting that, even for a relatively high porosity sandstone, the early  
443 deformation bands are dilatant except for low Lode angles.

444 For all studied loading paths, the continuous evolution of strain localiza-  
445 tion towards the end of the pre-peak regime leads to the formation of mature  
446 deformation bands, which develop concurrently to rapid stiffness weakening  
447 effects and a well defined transition of the mechanical response from a pre-  
448 peak to a post-peak regime. The development of these mature bands tends  
449 to occur in regions of the visible surface where early deformation bands have  
450 been persistent during the late pre-peak regime. After the peak stress, during  
451 increments 5-6 and 6-7 of Figure 7, most loading paths are seen to promote

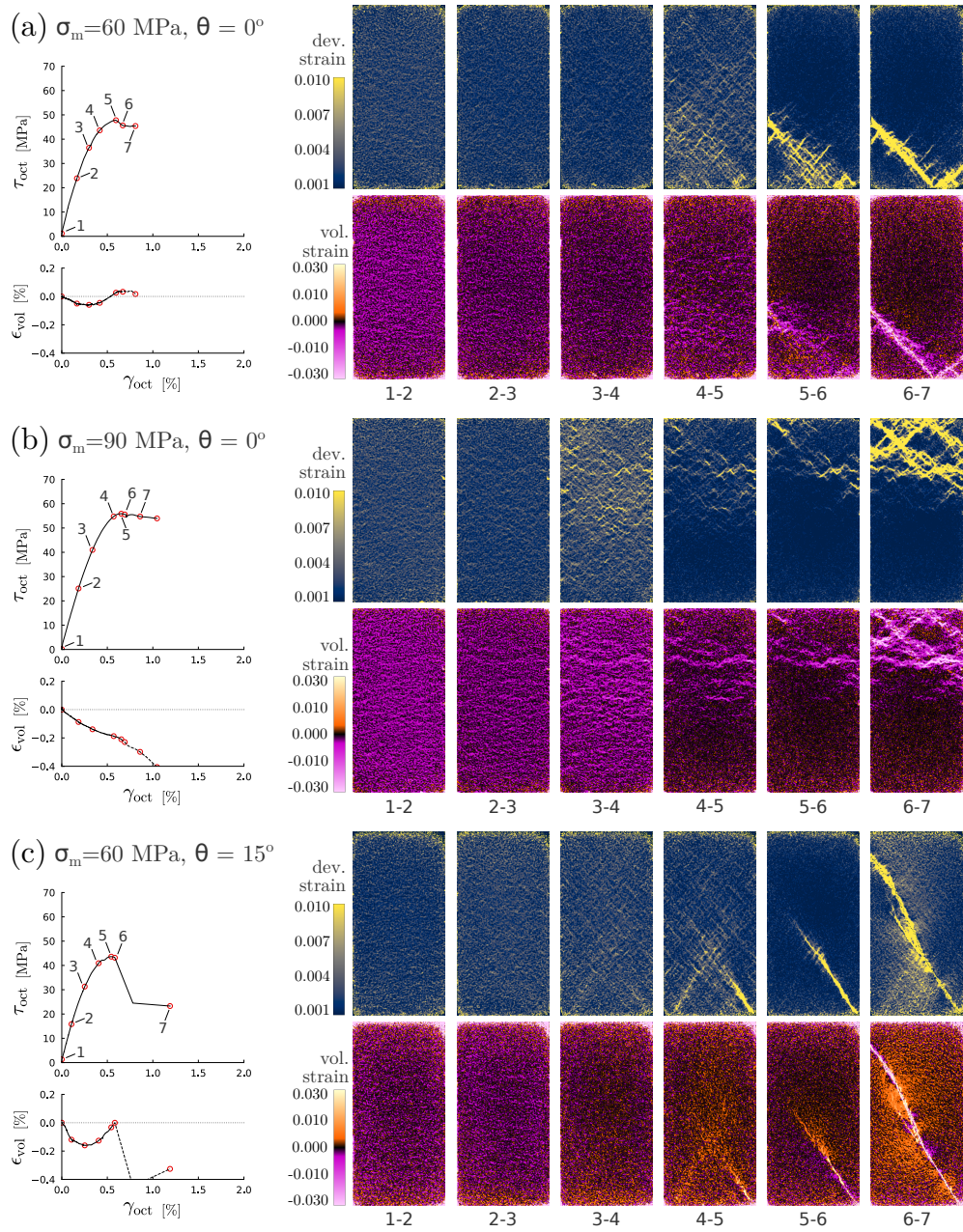


Figure 7: (part 1) Evolution of deviatoric and volumetric strain fields over the 2D ROI, for selected consecutive loading increments. The stress state and 3D volumetric deformation corresponding to each photograph are identified on the macroscopic curves by red circles.

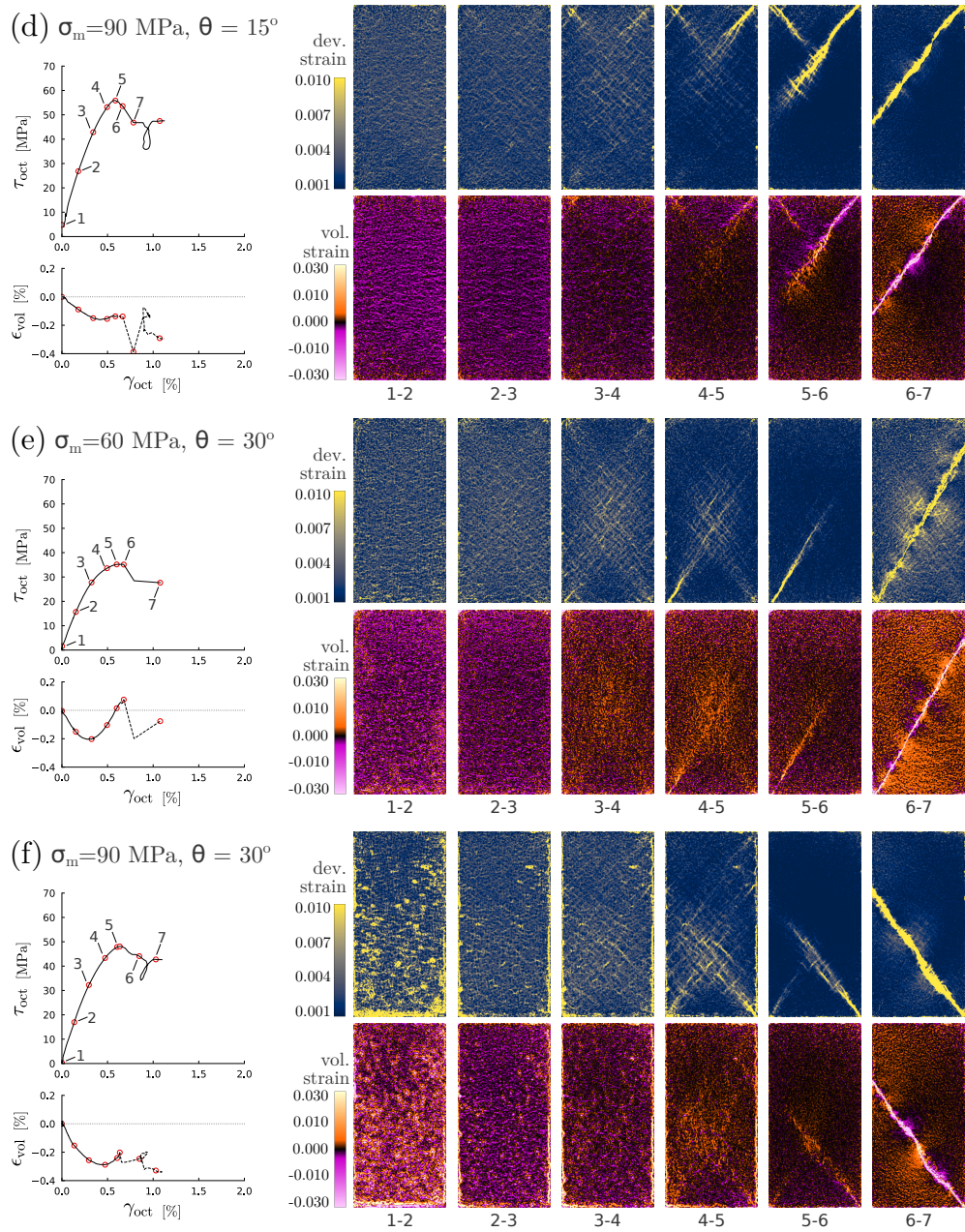


Figure 7: (part 2) Evolution of deviatoric and volumetric strain fields over the 2D ROI, for selected consecutive loading increments. The stress state and 3D volumetric deformation corresponding to each photograph are identified on the macroscopic curves by red circles.

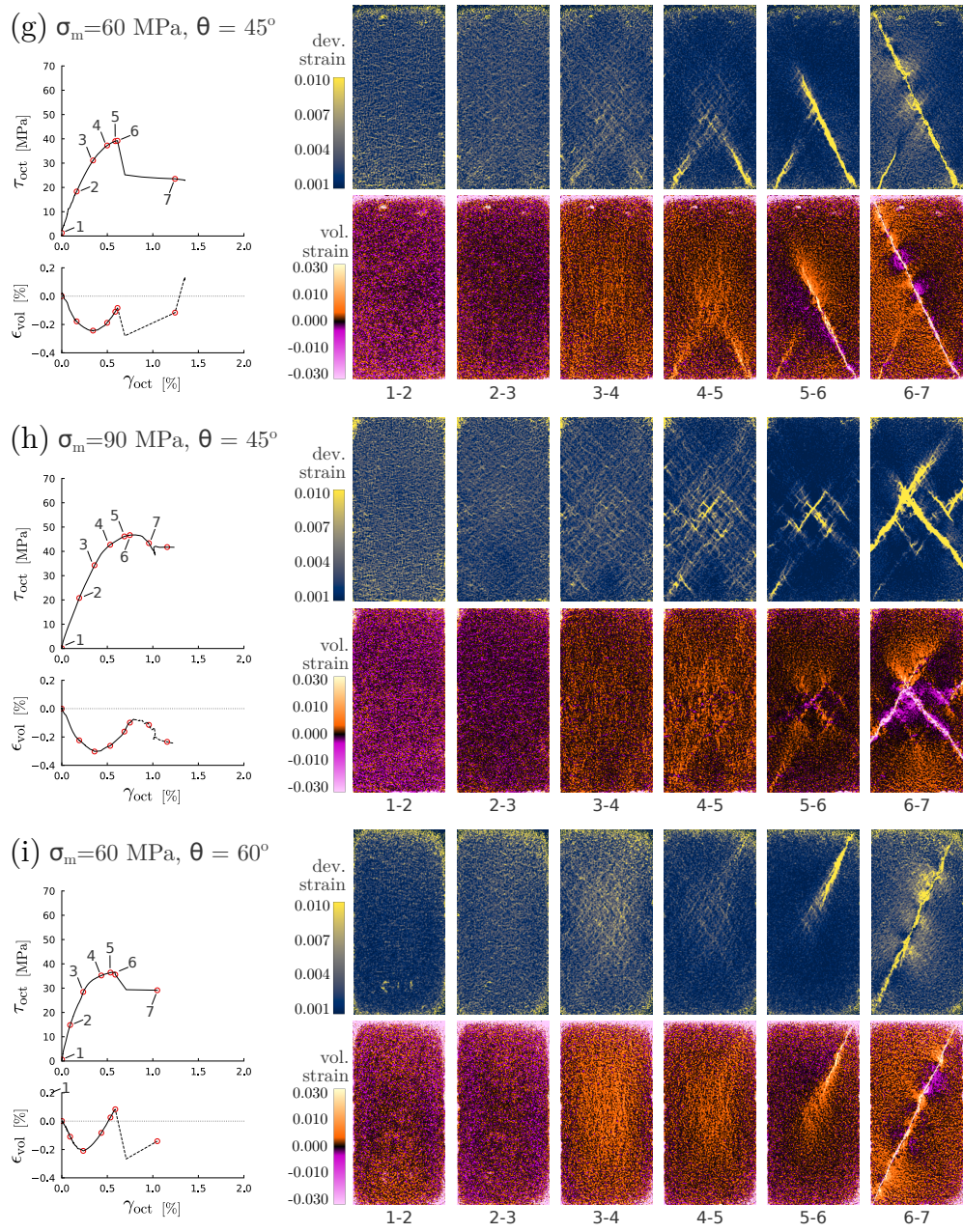


Figure 7: (part 3) Evolution of deviatoric and volumetric strain fields over the 2D ROI, for selected consecutive loading increments. The stress state and 3D volumetric deformation corresponding to each photograph are identified on the macroscopic curves by red circles.

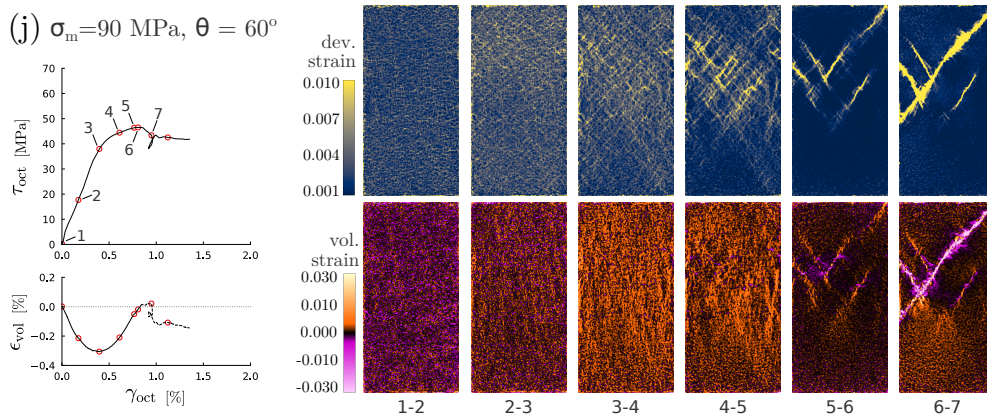


Figure 7: (part 4) Evolution of deviatoric and volumetric strain fields over the 2D ROI, for selected consecutive loading increments. The stress state and 3D volumetric deformation corresponding to each photograph are identified on the macroscopic curves by red circles.

452 a rapid concentration of strain localization into a unique linear structure  
 453 propagating through the sample. Once such a well developed localized re-  
 454 gion starts to form, it tends to attract additional deformation. A notable  
 455 exception to this deformation mode is identified for the loading path in Fig-  
 456 ure 7b (at  $\sigma_m = 90$  MPa and  $\theta = 0^\circ$ ), where mature deformation bands are  
 457 continuously activated and seen to relocate during the post-peak regime.

458 For samples tested at a mean stress of 60 MPa and a Lode angle greater  
 459 than  $0^\circ$ , the development of mature deformation bands ultimately results  
 460 in a global mechanical instability just after the peak stress, which can be  
 461 explained by the relatively low stiffness of the apparatus compared to the  
 462 absolute value of the tangential stiffness of the samples during the post-peak  
 463 regime. Under these conditions, the deformation band propagates quasi-  
 464 statically approximately halfway through the sample before quasi-static con-  
 465 ditions are lost. At this instant, a complete propagation of the localized zone  
 466 occurs dynamically, as elastic strain energy accumulated in both the sample  
 467 and the apparatus, is suddenly released. For this type of failure mechanism,  
 468 the resulting deformation zone is mostly concentrated around a unique and  
 469 narrow band, with a normal of the band oriented at a moderately high angle  
 470 ( $45^\circ$  to  $65^\circ$ ) with respect to the most compressive principal stress.

471 At the higher mean stress of 90 MPa and Lode angles greater than  $0^\circ$ ,  
 472 while mechanical instability can occur (see pressure regulation loops in Fig-  
 473 ure 5), the softening response is generally better controlled. At this higher

474 mean stress level, a number of localized structures evolving from early de-  
475 formation bands are more prone to coalesce and form a complex localized  
476 zone during the inception of a mature deformation band. The outcome of  
477 this coalescence process is a higher variability in the band thickness and in-  
478 ternal orientation, which is not entirely captured by a single characteristic  
479 length. In increment 6-7 of Figure 7, a combination of compactant and dila-  
480 tant zones is observed around the mature deformation band, corresponding  
481 to alternating localization orientation along the deformation band.

482 A substantial difference in the manifestation of mature deformation bands  
483 is observed, at both mean stresses, for a Lode angle of  $0^\circ$ . For these partic-  
484 ular ASC loading paths, it can be noted that the deformation bands are not  
485 necessarily polarized in the 1-3 visible plane since the minor and intermedi-  
486 ate principal stresses are equal. In the presented strain fields, the mature  
487 deformation bands evolve through the continuous intensification of extended  
488 localized regions close to the peak stress and later in the softening regime.  
489 This mode of localization occurs under a systematically compactant volu-  
490 metric strain of the mature deformation bands. Additionally, at the mean  
491 stress of 90 MPa, localized deformation regions evolve with the activation  
492 (and deactivation) of conjugated and parallel bands during the post-peak  
493 regime, resulting in a continuous relocation of the localized regions.

494 It is noted that in the majority of cases, but not in all, the mature bands  
495 pass through one of the corners of the sample. This corner attraction effect  
496 is well known in axisymmetric tests on cylindrical samples. However, the  
497 field measurements presented here clearly show that, probably thanks to the  
498 quality of the lubrication and the flatness of the surfaces, the deformations  
499 prior to the mature localization do not present any particular characteristics  
500 linked to the edges or corners of the samples. These appear to have a negli-  
501 gible impact on the overall response of the sample but sufficient to influence  
502 the position of the mature bands (not their orientation).

### 503 *3.3. Deformation band angles and dilatancy angle*

504 The orientation of early deformation bands ( $\beta^*$ ) and mature deforma-  
505 tion bands ( $\beta$ ) in the visible plane is measured as the angle between the  
506 normal vector to the quasi-linear localized regions and the direction of the  
507 most compressive principal stress ( $\sigma_1$ ). By characterising the deformation  
508 band angles during the loading phase from the incremental strain fields, the  
509 measure of the band angles reflects the intrinsic response of the material



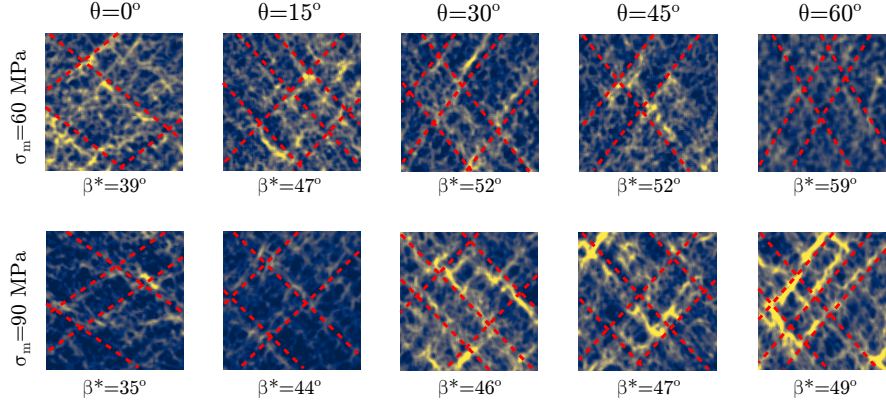


Figure 8: Identification of early deformation band angles ( $\beta^*$ ) from deviatoric strain field increments.

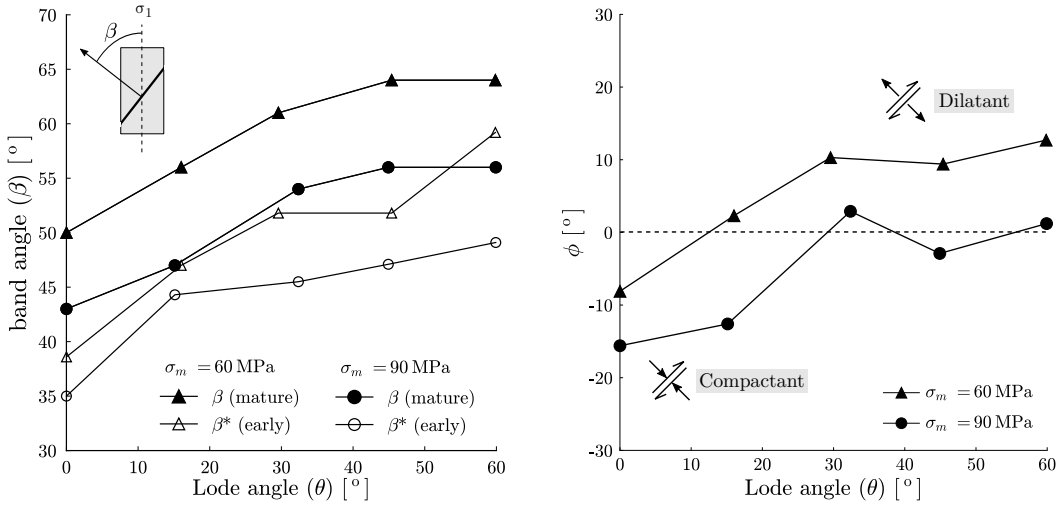


Figure 9: Evolution of the early and mature deformation band angles for the different loading paths in the octahedral plane(a), and dilatancy angle of mature deformation bands measured at their inception (b).

510 at the inception of the localized zones. Therefore, subsequent structural ef-  
 511 fects influencing the final band orientation, such as attraction to conjugated  
 512 bands or boundary imperfections during propagation, can be avoided. This  
 513 direct interpretation of the band angle at its inception is usually lost during  
 514 a post-mortem assessment of the failed material.

515 The predominant angle  $\beta^*$  characterizing early deformation bands is qual-

516 itatively assessed for the series of parallel and conjugated bands observed in  
517 the central sub-region of the ROI represented in Figure 8. While a certain  
518 variability in the early band angles in each sample can be observed, the an-  
519 gle of the red lines represents the general orientation of these bands as they  
520 emerge in samples subjected to the different loading paths.

521 The measured angle of early and mature deformation bands is reported for  
522 each experiment in Figure 9a. In both cases, the angle is seen to increase with  
523 decreasing mean stress and increasing Lode angle. At both mean stresses,  
524 the change in deformation band angle is generally more pronounced at lower  
525 Lode angles and tends to become constant at the higher Lode angles. The  
526 mature deformation bands are also systematically developing at a higher  
527 angle than the preceding early deformation bands, regardless of the loading  
528 path.

529 The dilatancy angle of mature deformation bands is defined as

$$\tan \phi = \frac{\Delta D^v}{\Delta D^s} , \quad (3)$$

530 where  $\Delta D^v$  and  $\Delta D^s$  are respectively the volumetric and shear offsets across  
531 the localized region [24]. With large measurement errors due to the speckle  
532 quality degradation inside the highly deformed bands, the volume and shear  
533 offsets are obtained from the averaged displacement on each side of the band  
534 during increment 5-6. In Figure 9b the dilatancy angle, positive for a dilatant  
535 band, is seen to increase with decreasing mean stress and increasing Lode an-  
536 gle, in a trend that is comparable to the evolution of the mature deformation  
537 band angle. Therefore, experiments at a low Lode angle show a shear local-  
538 ization with compaction (compacting shear band) and at a high Lode angle,  
539 a shear localization with dilation (dilatant shear band). Nevertheless, in this  
540 simplified representation, the complexities observed in the mature deforma-  
541 tion bands are not captured and must be kept in mind. This is particularly  
542 illustrated at the high mean stress of 90 MPa, where there are fluctuations  
543 in local orientations that impart a kind of roughness to the mature bands,  
544 as well as variability in volumetric strain along the band associated with this  
545 roughness.

#### 546 4. Discussion

547 Experimental results from this campaign on a high porosity sandstone  
548 provide strong evidences that distinct loading paths, following deviatoric

549 stress increments in the octahedral plane, promote substantially different  
 550 mechanical and kinematic responses. The influence of the mean stress on the  
 551 stress-strain behavior and failure mechanism, widely recognized in the case of  
 552 axisymmetric (compression and extension) stress paths, is here also observed  
 553 at intermediate values of the Lode angles. The effect of the different loading  
 554 paths is also evidenced by the local strain evolution measured from full-field  
 555 DIC. The significance of localization precursors, their evolution leading to  
 556 different failure mechanism and the transition in the brittle to ductile regime  
 557 is here further discussed.

#### 558 4.1. Localization Participation Ratio

559 To compare the rich information conveyed by the spatial distribution  
 560 of local strain measurements across different experiments, and quantify its  
 561 evolution during loading, the Localization Participation Ratio (LPR), first  
 562 introduced to the field of geomechanics by [25], is calculated for each series of  
 563 increments. This scalar index of localization provides an appreciation of the  
 564 deviatoric strain concentration over the visible plane. It is mathematically  
 565 defined as

$$\text{LPR} = \frac{1}{N_p} \frac{[\sum \epsilon_{dev}]^2}{\sum [\epsilon_{dev}^2]}, \quad (4)$$

566 where  $\epsilon_{dev}$  is the spatially distributed local deviatoric strain (as defined in  
 567 Appendix A), and  $N_p$  is the number of sampling points over the ROI. It is  
 568 calculated for each incremental field of deviatoric strain, quantifying a depart-  
 569 ure from a homogeneous spatial distribution. Therefore, calculated values  
 570 of the LPR are independent from the size of the selected increment, consider-  
 571 ing an equivalent homogeneous distribution of deviatoric deformation in this  
 572 increment is above the noise level resulting from the image acquisition and  
 573 treatment process. The upper limit,  $LPR = 1$ , corresponds to an homoge-  
 574 neous distribution of the considered spatial variable; whereas the lower limit  
 575 depends on the number of measurements on the domain for which a strongly  
 576 localized state (e.g. a Dirac distribution) results in  $LPR = \frac{1}{N_p}$ . The ROI  
 577 over the considered two dimensional plane is a  $600 \times 300$  area, excluding sub-  
 578 regions on the fringe of the sample surface where the autocorrelation radius  
 579 or DIC residuals were showing larger uncertainties.

580 Figure 10 shows the evolution of the LPR for the incremental deviatoric  
 581 strain fields reported in Figures 7. For all tested samples, the LPR indicates  
 582 a weak variation in diffuse localization during the pre-peak regime until a

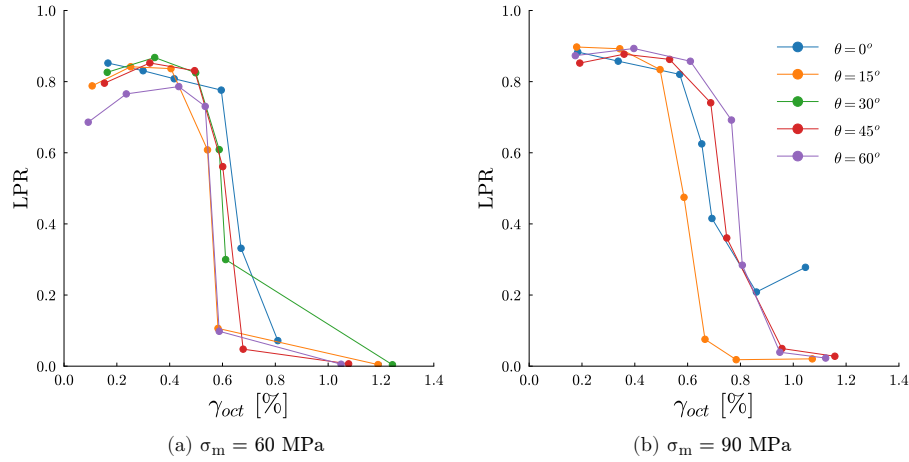


Figure 10: Localization Participation Ratio (LPR) evolution with octahedral strain for tests at (a)  $\sigma_m=60$  MPa and (b)  $\sigma_m=90$  MPa . The position of each point correspond to the end of the interval in which the LPR is measured. The sequence of increments are labeled in one of the curve as a reference.

583 strong decrease is initiated during the increment 4-5 (near the peak-stress),  
584 characteristic of well-developed early deformation band and the inception  
585 of a mature deformation band. At the mean stress of 60 MPa, the LPR  
586 decreases rapidly at around 0.6% octahedral strain. At a mean stress of  
587 90 MPa, for Lode angles of 0 and 15° the increase in the mean stress is  
588 seen to induce a decrease in the LPR at around the same octahedral strain  
589 of 0.6%. At higher Lode angles, the octahedral strain at which the LPR  
590 starts to decrease is however noticeably higher and shows a more progressive  
591 evolution towards the full development of mature deformation bands. Even if  
592 the spatial organization of deviatoric strain is not captured in the LPR, this  
593 observation is consistent with the progressive evolution in the coalescence  
594 of localized deformation at higher mean stress and for Lode angles above  
595 15°. This dissimilar kinematic response to an increase in the mean stress,  
596 between Lode angles above and below 30°, suggests a strong influence of the  
597 Lode angle on the deformation regime leading to failure.

#### 598 4.2. From localization precursors to failure

599 During the early deviatoric loading phase, corresponding to increments  
600 1-2 and 2-3, the relatively large inelastic compaction observed at the sam-  
601 ple scale for experiments conducted at higher Lode angles occurs in a rela-  
602 tively diffuse deformation process, which is reflected by the initially stable or

603 slightly increasing LPR. Therefore, during this phase of loading, the sample  
604 deforms mostly in the absence of identifiable structural localization patterns.  
605 Nonetheless, samples experiencing a higher initial compaction at the sample  
606 scale, are associated with a subsequent well defined transition into a dilatant  
607 regime, during increment 3-4. Concurrently, these samples present a higher  
608 deviatoric strain intensity and dilatant volumetric strains concentrated inside  
609 early deformation bands. This relation between the initial compactant be-  
610 havior and the later localization behavior at different Lode angle emphasizes  
611 the role of selected loading paths in a single octahedral plane in promoting  
612 different diffuse deformation responses and transitions into the development  
613 of localization structures.

614 The emergence of dilatant early deformation bands results from the pro-  
615 gressive clustering and densification of initially scattered localization precu-  
616 sors. Under monotonic deviatoric loading conditions in the octahedral plane,  
617 the inception and intensification of these early bands appear to be driven by  
618 shear deformation with a predominantly dilatant behavior. In this instance,  
619 the dilatant volumetric strain, mostly apparent in localized zones at higher  
620 Lode angles, is concurrent with a reversal from an incremental compactancy  
621 to an incremental dilatancy in the global volumetric deformation measured  
622 at the sample scale. This observation suggests that the presence of dilatant  
623 early deformation bands constitute an important deformation mechanism  
624 contributing to inelastic deformation and a concentration of weakening defor-  
625 mation in fewer regions of the porous rock during the later pre-peak regime.  
626 The continued development of early deformation bands appears to be an im-  
627 portant phenomenon concomitant with the loss of linearity in the octahedral  
628 stress-strain curve. It can therefore be inferred that, at least for the range of  
629 mean stress investigated here, the lower octahedral stress attained at higher  
630 Lode angles is related to the emergence and continuous organization of these  
631 early structural deformation zones developing well before the peak-stress.

632 In samples where a well defined dilatant regime is observed, i.e. for load-  
633 ing paths where  $\theta > 0^\circ$ , strain localization tends to rapidly coalesce and  
634 concentrate into fewer localized regions of higher intensity and increasing di-  
635 latancy as the peak stress is approached. Consequently, during this process,  
636 previously localized regions tend to progressively become inactive. The peak  
637 stress is reached when a dominant sublinear zone of high strain concentra-  
638 tion propagates through the sample and develops into a mature deformation  
639 band by connecting early deformation bands of limited length into a contin-  
640 uous structure. During this transition, a dilatant region, comprising earlier

641 deformed regions, remains active ahead of the tip of the propagating ma-  
 642 ture deformation band (identified in Figure 11b-c). Since early deformation  
 643 bands tend to initially form at a lower angle than the later mature deforma-  
 644 tion bands, the localized zones become connected by sections of alternating  
 645 internal angles. Figure 11 illustrates the presence of echelon type patterns  
 646 connecting early deformation bands of higher activity along the length of  
 647 the mature localized region, with associated dilatant (high angle) and com-  
 648 pactant (low angle) zones. The prevalence of this deformation mechanism  
 649 is corroborated by similar, although post mortem, observations in [7]. In a  
 650 series of experiments on Dunham dolomite in true triaxial loading conditions,  
 651 the author identifies the existence of secondary micro cracks, here identified  
 652 as early deformation bands, on the sample surface, as well as a variation in  
 653 the internal angles of mature deformation bands along its length. Conse-  
 654 quently, the formation of early deformation bands, here identified at a lower  
 655 inclination, seems to play a significant role on the roughness and thick-  
 656 ness variation of mature deformation bands, which in turns can influence  
 657 the post-failure behavior of the material, during the quasi-stable permanent  
 658 regime. This suggests that the spatially inhomogeneous deformation and in-  
 659 duced anisotropy from oriented early localized structures can influence the  
 660 later occurring failure mode in the material.

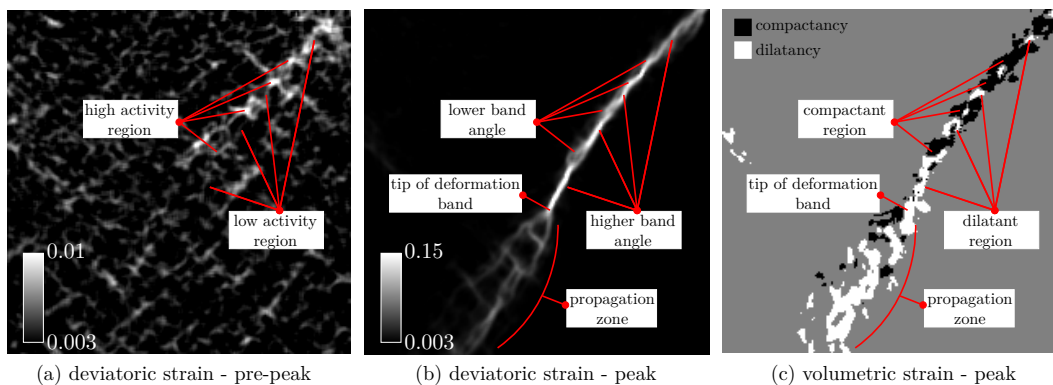


Figure 11: Correspondence between early localized regions (a) during the pre-peak and the mature deformation band angle variation (b) and volumetric strain (c) at the peak. The represented sub-regions are increments 3-4 (pre-peak) and 5-6 (peak) of Figure 7d.

661 While compactant sub-regions along propagating deformation bands are  
 662 observed at higher Lode angles, a contrasting continuously compactant type  
 663 of deformation bands is only observed in samples tested at a Lode angle of

664 0°. For these particular ASC loading paths, deformation mechanisms differ  
665 significantly in terms of the development of localization structures during the  
666 transition around the peak stress. In this case, strain localization appears  
667 to be more distributed across the sample and develops simultaneously in  
668 multiple and extended active regions. During the late pre-peak to post-  
669 peak transition, the distribution of compactant deformation bands over a  
670 larger part of the sample suggests that, unlike dilating deformation bands,  
671 compactant bands have a lesser tendency to concentrate strain localization  
672 into a single linear mature deformation band, as previously suggested by  
673 [18]. Instead, the stable deformation process inside compactant band result  
674 in a network of evolving and highly connected internal sub-structures. This  
675 phenomenon is most apparent at the higher mean stress of 90 MPa, where  
676 multiple localized regions remain active during the post-peak regime and  
677 progressively relocate to nearby lesser deformed regions of the sample.

#### 678 *4.3. Brittle-ductile transition*

679 For the range of stress paths studied in this experimental campaign, the  
680 variation in the mean stress and Lode angle between different loading paths  
681 is shown to have a significant influence on the global mechanical response  
682 and on the spatial evolution of localization phenomena. For the studied  
683 Vosges sandstone, an increase in the mean stress or a decrease in the Lode  
684 angle systematically induce a transition from a brittle to a ductile mode of  
685 failure. This transition is apparent in the mechanical response by an increase  
686 in the maximum octahedral stress, a progressive softening in the post peak  
687 and a reduction in the incremental volumetric strain at the peak. In terms of  
688 localized deformation modes, the brittle to ductile transition is seen through a  
689 decrease in both the band angle of early and mature deformation bands and in  
690 the dilatancy angle measured at the peak stress. Furthermore, concentrated  
691 dilatant early and mature deformation bands are systematically observed in  
692 samples showing a brittle failure mode, while a compactant behavior and  
693 relocation of mature deformation bands is associated with a ductile failure  
694 mode.

695 The transition in the brittle-ductile behavior in the deviatoric plane is  
696 further evidenced by the non-linear shift in the peak octahedral stress oc-  
697 ccurring in the low to intermediate interval of Lode angles (Figure 6). The  
698 similarity in the peak value between 0° and 15° at the mean stress of 90 MPa,  
699 reinforced by an additional point at 15° (Table 2), emphasizes the observa-  
700 tion of a shorter transition range at higher mean stress. Past this transition,

701 further increase of the Lode angles appears to have a diminishing effect on  
702 the peak stress and mature deformation band angles.

703 In terms of strain localization patterns, a brittle-ductile transition can  
704 generally be associated with the aforementioned change in the type of deforma-  
705 tion bands formed during the post-peak regime. A shift in the deformation  
706 patterns towards a ductile response is seen to occur with a decrease in the  
707 Lode angles. In a study following similar loading paths at imposed mean  
708 stress and Lode angle, Ingraham et al. [13] observed a similar trend, where  
709 a decrease in the Lode angle results in a shift from well defined narrow de-  
710 formation bands to a more diffuse mode of failure, where individual bands  
711 could not be identified in a post-mortem assessment. The authors also noted  
712 that the shift towards a more ductile response occurs at higher mean stresses  
713 for higher Lode angles.

714 Despite a general concordance between the mechanical and kinematic (i.e.  
715 localization pattern) interpretation of the brittle-ductile transition, it can be  
716 noted that, for samples tested at a mean stress of 90 MPa and Lode angles of  
717  $0^\circ$  and  $15^\circ$ , while having a similar octahedral stress-strain response during the  
718 pre-peak regime and a close peak octahedral stress, their mature localization  
719 patterns is quit different. At the Lode angle of  $0^\circ$ , a ductile (continuously  
720 relocating and compactant) deformation mode is observed, while at  $15^\circ$  the  
721 deformation band can be seen as more brittle (propagating and compactant  
722 with dilatant zones). Therefore, in this case the ductility in the deformation  
723 mode is not uniquely related to the macroscopic stress-strain response, but  
724 is specific to the loading path in the octahedral plane.

725 The range of the brittle-ductile transition is concurrent with a progres-  
726 sive shift in the intermediate principal strain (direction 2) with increasing  
727 Lode angle and the increasing dilatant behavior observed in early deforma-  
728 tion bands. From these experimental observations, it appears that stress  
729 paths promoting a shortening of the sample in the intermediate principal  
730 direction 2 is contributing to localized dilatant effects in the visible plane  
731 (1-3). The concentration of this effect in early deformation bands implies  
732 that this mechanism probably exacerbates the alteration and weakening of  
733 the rock fabric in localized zones.

734 In loading paths promoting a transition towards a ductile failure, the  
735 total octahedral strain in the sample at the inception of the softening regime  
736 is equivalent or lower than for loading paths leading to a more brittle failure.  
737 Therefore, the higher peak octahedral stress attained in those experiments  
738 seems to be attributed to the slower bending in the stress strain response,



739 rather than a capacity of the samples to sustain higher global deformation,  
740 as the mechanical response transitions toward the ductile regime.

## 741 **5. Concluding remarks**

742 The presented experimental study, investigating the mechanical behavior  
743 and full-field kinematics of a high porosity sandstone, contributes original  
744 results obtained under unconventional and rarely studied true triaxial loading  
745 paths in the octahedral plane. Using a unique experimental setup, strain  
746 localization structures and their evolution during the loading phase were  
747 assessed and compared for different deformation regimes. The use of optical  
748 images, acquired at a high spatial and temporal resolution, and analysed  
749 through digital image correlation (DIC), proved to be an effective method to  
750 explore the role of micro-kinematics on the inception of localization structures  
751 leading to different failure modes.

752 This study highlighted important transitions in the material response  
753 between the series of experiments performed at different purely deviatoric  
754 loading paths. Notably, it is demonstrated that an increase in the prescribed  
755 Lode angle results in a systematic increase in the angle of early (pre-peak)  
756 and mature (peak and post-peak) deformation bands, an increase in the  
757 dilatancy angle of mature deformation bands, as well as a decrease in the  
758 peak octahedral stress and an embrittlement of the sample response.

759 The evolution of kinematic structures identified at the surface of the sam-  
760 ple during the pre-peak regime clearly showed the emergence and continuous  
761 development of early deformation bands. These bands are seen to organize  
762 into parallel and conjugated sublinear strain localization regions of limited  
763 length well before the peak stress. The ubiquitous presence of this inho-  
764 mogeneous mode of deformation seems to play an important role in the de-  
765 velopment of inelastic and dilatant effects at the sample scale. Also, the  
766 observation of early strain localization is concordant with the continuous  
767 weakening of the rock fabric integrity during the pre-peak regime. During  
768 this loading phase, the progressive intensification and concentration of strain  
769 in early deformation bands culminate into the formation of either dilatant  
770 mature deformation bands propagating through the sample (fragile) or the  
771 continuous deactivation/relocation of larger localized regions during the soft-  
772 ening regime (ductile). For the range of mean stresses investigated in this  
773 study, the brittle-ductile transition, observed in both the strain localization

774 patterns and the mechanical behavior, occurs mostly in the lower to inter-  
 775 mediate range of Lode angles, between  $\theta = 0^\circ$  and  $30^\circ$ .

776 Finally, the angle of early deformation bands is found to be lower than  
 777 the angle of mature deformation bands. Therefore, as early localized regions  
 778 become connected, the alternating variation in internal angle of mature de-  
 779 formation bands influences their aspects (departure from flatness) and thus  
 780 the material residual strength in the permanent regime.

781 These observations are important in taking into account the non-linear  
 782 and complex response of porous rocks subjected to general true triaxial load-  
 783 ing paths. Therefore, the relation between the mechanical response and  
 784 development of strain localization, especially in the early stages of loading,  
 785 should be considered to enrich existing constitutive models for high poros-  
 786 ity sandstones, as well as other analog materials characterized by cohesive-  
 787 granular micro-structures.

## 788 Appendix A. Strain decomposition

789 In the Software for Practical Analysis of Material, the deformation gra-  
 790 dient tensor,  $\mathbf{F}$ , is derived from the incremental displacement field obtained  
 791 by 2D digital image correlation. In the large scale framework implemented  
 792 in the strain calculation code, the change in length and angle is represented  
 793 by the right stretch tensor:

$$\mathbf{U} = \sqrt{\mathbf{F}^T \cdot \mathbf{F}} = \sqrt{\mathbf{U}^T \cdot \mathbf{R}^T \cdot \mathbf{R} \cdot \mathbf{U}} = \sqrt{\mathbf{U}^T \cdot \mathbf{U}} \quad (\text{A.1})$$

794 where  $\mathbf{R}$  is the antisymmetric part of  $\mathbf{F}$ , i.e. the rotation matrix, and " $\cdot$ "  
 795 denotes the inner product operator. This tensor is further decomposed into  
 796 its isotropic and deviatoric multiplicative parts,

$$\mathbf{U} = \mathbf{U}_{iso} \cdot \mathbf{U}_{dev} \quad (\text{A.2})$$

797 with

$$\mathbf{U}_{iso} = J^{1/2} \cdot \mathbf{I} \quad (\text{A.3})$$

798 and

$$\mathbf{U}_{dev} = \frac{1}{J^{1/2}} \cdot \mathbf{U} \quad (\text{A.4})$$

799 where  $J$  is the Jacobian, the determinant of  $\mathbf{F}$ , giving the change in volume  
 800 between the initial and final configurations. Therefore, the two scalar quanti-  
 801 ties representing the volumetric and deviatoric deformations are respectively,

$$\epsilon_{vol} = J - 1 \quad (\text{A.5})$$

802 and the Euclidean norm of deviatoric part of the stretch tensor,

$$\epsilon_{dev} = \|\mathbf{U}_{dev}\| \quad (\text{A.6})$$

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