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Manuscript accepted by *Geochemistry, Geophysics, Geosystems*

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Key Points:

- Time variations of a noise-derived *P*-type phase are compared with those of the ocean wave heights and microseism sources.
- Do not equate a positive correlation with a causal relation when studying the links between noise sources and noise-derived signals.
- The derivation of seismic signals from ambient noise relies on the competition between the effective and ineffective sources.

Abstract

Seismic signals can be extracted from ambient noise wavefields by the correlation technique. Recently, a prominent *P*-type phase was observed from teleseismic noise correlations in the secondary microseism period band. The phase is named P_{dmc} in this paper, corresponding to its origin from the interference between the Direct *P* waves transmitting through the deep Mantle and the Core (*P* and *PKPab* waves). We extract the phase by correlating noise records from two seismic networks in the northern hemisphere, and locate the microseism sources that are efficient for the P_{dmc} construction in the south Pacific. We investigate the spatiotemporal links of the P_{dmc} signal with global oceanic waves and microseism sources. Interestingly, the correlation with wave height is higher in several regions surrounding the effective source region, rather than in the effective source region. The P_{dmc} amplitude is highly correlated with the power of the effective microseism sources. Also, it is apparently correlated with ineffective sources in the southern hemisphere, and anti-correlated with sources in the northern hemisphere. We ascribe the correlation with the ineffective southern sources to the spatiotemporal interconnections of the southern sources. The anti-correlation with northern sources can be explained by the reverse seasonal patterns between the southern and northern sources, and by that the northern sources impede the signal construction. The signal construction from noise correlations relies on the

competition between the effective and ineffective sources, not just on the power of the effective sources. This principle should be valid in a general sense for noise-derived signals.

Plain Language Summary

Earth is experiencing tiny but incessant movement induced by natural forces, particularly, storm-driven ocean waves. While this ambient seismic noise (microseism) was deemed a nuisance in the past, it can be turned into signals via the seismic correlation technique.

Recently, a new *P*-type phase was derived from the noise correlations between two regional seismic networks. The noise-derived phase originates from the correlation between *P* waves that propagate through the deep mantle and outer core of the Earth.

The temporal amplitude variations of the noise-derived signals are compared with the variations of microseism sources in the oceans. We show that the signal emergence depends on the competition between the sources in a specific region that contribute to the signals and sources in other regions. The conclusion can be generalized to other noise-derived seismic phases.

We also analyze the links of the noise-derived signals to ocean waves. In our case, the ocean waves in the contributing source region are dominated by wind seas forced by local winds, whereas the excitation of microseisms is primarily owing to the freely traveling swells generated by oceanic storms in surrounding regions.

1 Introduction

The incessant background vibrations of Earth had been observed as early as the birth of seismometers in the later 19th century (Bernard, 1990; Dewey & Byerly, 1969; Ebeling, 2012). They were termed "microseisms" due to their tiny amplitudes. With more apparatus deployed worldwide, it was soon recognized that microseisms are ubiquitous and irrelevant to seismicity. The observation of microseisms aroused interests from various disciplines. Researchers linked the generation of microseisms to atmosphere processes and ocean wave activity. Meteorologists tried to employ land observations of microseisms to track remote oceanic storms (e.g., Harrison, 1924). Since the mid-twentieth century, it has been well known that microseisms are excited by storm-driven ocean waves. The most energetic microseisms that dominate the seismic noise spectra, namely, the so-called secondary microseisms at seismic periods around 7 s (Peterson, 1993), are excited by the nonlinear interactions between nearly equal-frequency ocean waves propagating in nearly opposite directions (Longuet-Higgins, 1950; Hasselmann, 1963). The periods of the excited secondary microseisms are half those of the colliding ocean waves. The excitation source is equivalent to a vertical pressure applied to the water surface, which is proportional to the product of the heights of the opposing equal-frequency waves. Due to this second order relation, moderate sea states can sometimes generate loud microseism noise (Obrebski et al., 2012). Thus, the presence of a strong microseism event does not necessarily imply a locally intense sea state.

By coupling the excitation theory of secondary microseisms proposed by Longuet-Higgins (1950) with the ocean wave action model, Kedar et al. (2008) modeled the secondary microseism excitations in the north Atlantic, and validated the numerical modeling by comparing with inland seismological observations. Afterwards, more authors simulated the oceanic microseism sources and some reported the consistency between predictions and observations (e.g., Ardhuin et al., 2011, 2015; Hillers et al., 2012; Stutzmann et al., 2012; Nishida & Takagi, 2016). Stopa et al. (2019) compared the microseism simulations with real observations to validate their corrections to the global reanalysis wind fields, which systematically reduced the residuals in the wave hindcast over the past decades.

The seismic excitation by an oceanic microseism source is essentially akin to that by an earthquake, in that the seismic wavefield recorded at any point is a convolution of the source time function with the Green function of the propagating medium between source and receiver. Their main difference lies in the source process. For earthquakes, the sudden rupture of faults leads to short-duration, impulsive source time functions. Isolated seismic phases are generally distinguishable from the seismograms. In contrast, the excitation of microseisms, approximated as Gaussian random process by some authors (Peterson, 1993; Steim, 2015), is incessant, leading to long, random-like source time functions. The convolution mixture signals are not directly discernible from the seismograms. With array beamforming (Rost & Thomas, 2002) or correlation technique (Campillo & Paul, 2003; Shapiro & Campillo, 2004), specific phases from distant microseism sources have been identified from microseism noise records (e.g., Gerstoft et al., 2008; Landès et al., 2010; Zhang et al., 2010; Euler et al., 2014; Reading et al., 2014; Gal et al., 2015; Liu et al., 2016; Nishida & Takagi, 2016; Meschede et al., 2017, 2018; Retailleau & Gualtieri, 2019). The correlation technique is advantageous in that, by correlating the noise records at two receivers, explicit seismic signals can be derived. Noise-derived surface waves have been used to infer the azimuthal and seasonal changes of noise sources (e.g., Stehly et al., 2006). Noise-derived body waves can provide better constrains in imaging the noise sources (Landès et al., 2010). Recently, deep body waves that propagate through the mantle and core have been extracted from ambient noise (e.g., Boué et al., 2013; Lin et al., 2013; Nishida, 2013; Poli et al., 2015; Xia et al., 2016; Spica et al., 2017; Retailleau et al., 2020). The noise-derived body waves are valuable for surveying the deep structure and for understanding the links between seismological observations and atmospheric/oceanographic phenomena.

Hillers et al. (2012) made the first global-scale comparison between the oceanic microseism sources derived from seismological observations and oceanographic modeling. The seismologically derived data (time resolution: 13 days; spatial resolution: 2.5° latitude × 5° longitude) are the global back-projections of near-zero-lag *P* signals generated from the cross correlations of microseism *P* waves at seismic array (Landès et al., 2010). The modeled data (time resolution: 3 hours; spatial resolution: 1° latitude × 1.25° longitude) are a global extension of the numerical simulation by Kedar et al. (2008). The two datasets are resampled to common resolutions for comparison. For the seismologically derived data, the back-projection is based on the relationship between the source-receiver distance and the horizontal slowness of teleseismic *P* wave. However, seismic phases that have common slownesses (e.g., *P* and *PP* waves) cannot be discriminated in this method (Gerstoft et al., 2008; Landès et al., 2010). Thus, the imaged sources are somewhat ambiguous. For the modeled data, coastal reflections of ocean waves, that can play a role in the ocean wave-wave interactions at near-coast regions (Longuet-Higgins, 1950; Ardhuin et al., 2011), are neglected. Due to the resonance of seismic waves in the water columns, bathymetry can have significant effect on the excitation of microseisms (Longuet-Higgins, 1950; Kedar et al., 2008; Hillers et al., 2012). The importance to account for the bathymetric effect on the microseism *P*-wave excitations has been addressed in several studies (e.g., Euler et al., 2014; Gal et al., 2015; Meschede et al., 2017). Hillers et al. (2012) considered the bathymetric effect, but using the amplification factors derived by Longuet-Higgins (1950) for surface waves.

Rascle and Ardhuin (2013) established an oceanographic hindcast database that includes global oceanic secondary microseism sources of a 3-hour time resolution and a 0.5° spatial resolution. Coastal reflections were accounted for in the modeling (Ardhuin et al., 2011). Regarding the bathymetric effect on microseism excitations, Gualtieri et al. (2014) proposed the formulae for body waves based on ray theory. Concerning the localization of noise sources, Li et al. (2020) developed a double-array method that can estimate the respective slownesses of the interfering waves, and thereby, provide better constrains for the determination of the correlated seismic phases. The microseism sources that are effective for the derivation of seismic signals from noise records, can be mapped by back-projecting the noise-derived signals along the ray paths of the correlated phases. The double-array configuration eliminates the ambiguity in determining the effective source region (Fresnel zone). In this study, we integrate these new progresses to survey the associations of noise-derived body waves to ocean wave activity and microseism excitations.

This paper is organized as follows. In section 2, we review the main results of Li et al. (2020) who reported the observation of a prominent *P*-type phase from the noise correlations between two regional seismic networks at teleseismic distance. The noise-derived phase has its spectral content concentrated in the period band of the secondary microseisms that are excited by the nonlinear ocean wave-wave interactions. In this paper, we denote the phase as P_{dmc}, corresponding to the fact that the phase originates from the correlation between the Direct *P* waves that transmit through the deep Mantle and the outer Core (microseism *P* and *PKPab* waves). In section 3, we estimate the temporal variations in the P_{dmc} amplitude and refute the associations to seismicity. In section 4, correlation analysis is used to unveil the spatiotemporal links of the P_{dmc} signal with the global oceanic wave climate and microseism sources. Last, we discuss the significance of this study in seismology, oceanography and climate science.

2 Noise-derived P_{dmc} phase

Li et al. (2020) correlated the seismic noise records from two regional seismic networks at teleseismic distance: the FNET array in Japan and the LAPNET array in Finland (Fig. 1a). The continuous seismograms were divided into 4 h segments and whitened in the frequency domain. Segments with large spikes (like earthquakes) were discarded. The available segments of each FNET-LPANET station pairs were correlated. For more technical details, see section 2 of Li et al. (2020). From the vertical-vertical components of the FNET-LAPNET noise correlations, they observed coherent spurious arrivals (the P_{dmc} phase named in the previous section) that emerged ~200 s earlier than the direct *P* waves (Fig. 1b). By estimating the respective slownesses of the interfering waves and their time delay, it is unveiled that a quasi-stationary phase interference between the teleseismic *P* waves at FNET and the *PKPab* waves at LAPNET, emanating from noise sources in the ocean south of New Zealand (NZ), lead to the noise-derived P_{dmc} phase (Fig. 1c). The quasi-stationary phase condition refers to that the interfering waves have no common path or common slowness, but the stack of correlation functions over a range of sources can still be constructive as an effect of finite frequency. This observation contrasts with the strict stationary phase condition that has been employed by Phạm et al. (2018) to explain the spurious body phases in the earthquake coda correlations. The strict condition implies the existence of sources in the stationary-phase region, or say, the correlated waves have common ray paths or common slownesses. Li et al. (2020) substantiates the explanation of quasi-stationary phase for the observed P_{dmc} signals with numerical experiments based on ray theory and based on spectral-element modeling, and highlighted the discrepancies between (microseism) noise correlations and coda correlations.

The P_{dmc} phase has an apparent slowness of 4.6 s/deg, while the slownesses of the interfering *P* and *PKPab* waves are 4.7 s/deg and 4.2 s/deg, respectively. The dominant period of the P_{dmc} phase is 6.2 s, typical for secondary microseisms. The observation of the P_{dmc} phase is time-asymmetric (Fig. S1a). Its absence from the mirror side is ascribed to the faintness of the corresponding source in the low-latitude Atlantic (Fig. S1b).

There are several advantages to investigating the links between noise-derived signals and microseism sources with the P_{dmc} phase. First, the correlated *P* and *PKPab* waves are both prominent phases in the ballistic microseism wavefields. The P_{dmc} phase is easily observable from noise correlations, even between some single station pairs and on some single days (Fig. S2). Second, the isolation of P_{dmc} signals avoids potential bias caused by other prominent signals. Third, the effective sources are confined in a limited, unique region (Fresnel zone). In contrast, noise-derived surface waves have a broad Fresnel zone around the line across the correlated stations, and noise-derived *P* waves can have multiple Fresnel zones (see fig. 5 of Boué et al., 2014 for instance). The uniqueness of the effective source region can facilitate the study on the correlation between the noise-derived signals and the effective sources. Fourth, the correlated FNET and LAPNET networks are next to the northern Pacific and Atlantic, respectively, while the effective source region locates in the southern Pacific. The northern oceans have consistent seasonal variation pattern distinct from (reverse to) that of the southern oceans (Stehly et al., 2006; Stutzmann et al., 2009; Landès et al., 2010; Hillers et al., 2012; Reading et al., 2014; Turners et al., 2020). These geographical configurations make the observations easier to interpret. Last, there happens to be a seismic array (GEONET) in NZ next to the effective source region for the P_{dmc} phase. The seismic data from GEONET provide extra support to our study.

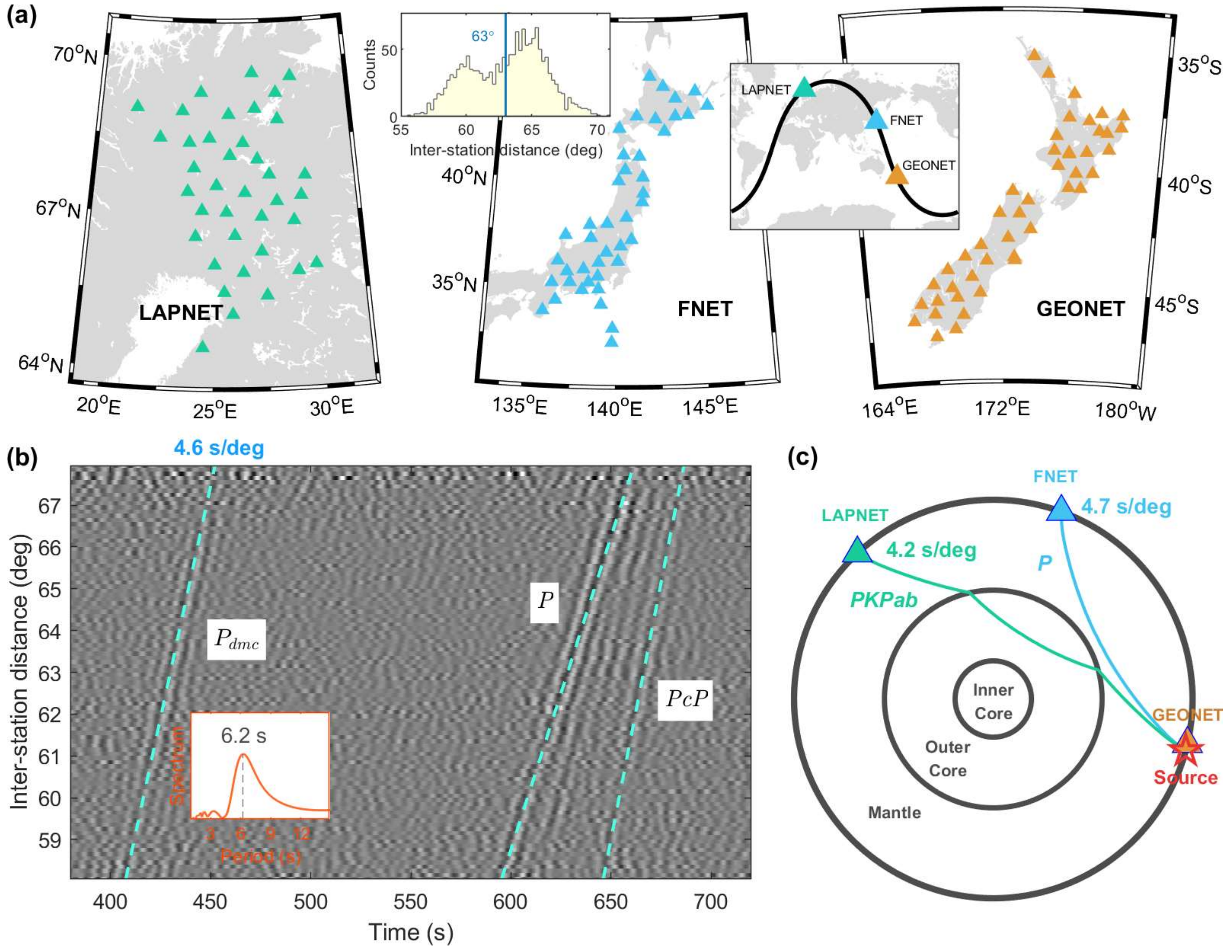

Figure 1. (a) Three regional broadband seismic networks used in this study: left, the LAPNET array in Finland (38 stations); center, the FNET array in Japan (41 stations); right, the GEONET array in New Zealand (46 stations). The histogram inset shows the distribution of the separation distances between the 1558 FNET-LAPNET station pairs. The center-to-center distance is 63° between LAPNET and FNET, and 85° between FNET and GEONET. The global inset shows the geographical locations of the three networks that are aligned on a great circle (dark line). (b) Annual FNET-LAPNET noise correlations that are filtered between 5 s and 10 s and stacked over time and in 0.1° inter-station distance bins. The spectrum inset indicates that the P_{dmc} phase has a 6.2 s peak period. (c) Ray paths of the interfering waves that generate the P_{dmc} phase. The effective source region is close to GEONET.

3 Temporal variations

We extract the temporal variations of the P_{dmc} signals by beamforming the FNET-LAPNET noise correlations on a daily basis. The daily noise correlations are shifted and stacked by

$$B(t) = \langle C_{ij}\big(t + (d_{ij} - d_0)\cdot p\big)\rangle, \qquad (1)$$

with $\langle\cdot\rangle$ the mean operator, C_{ij} and d_{ij} the correlation function and the distance between the ith FNET station and the jth LAPNET station, d_0 the reference distance (63°), p the apparent slowness of the P_{dmc} phase (4.6 s/deg), and t the time. The image in Fig. 2 shows the envelopes of the daily beams computed from the Hilbert transform of Eq. (1), with the daily P_{dmc} strength by averaging the envelope amplitudes plotted in the top panel. The strength of daily P_{dmc} signals varies strikingly, extremely strong on some single days (see the labeled dates in the P_{dmc} strength curve for examples), but indiscernible on most other days.

Considering that the region of effective source is tectonically active, one should investigate the plausible connection between the P_{dmc} signals and seismicity. From Fig. 2, it is obvious that P_{dmc} is decorrelated with the NZ seismicity. Also, it shows no connection with global large earthquakes as has been observed for coda-derived core phases at periods of 20 to 50 s (Lin & Tsai, 2013; Boué et al., 2014). That again demonstrates the substantial difference between ambient noise correlations and earthquake coda correlations, as emphasized by Li et al. (2020). The P_{dmc} strength exhibits an obvious pattern of seasonal variation. The seasonal pattern does not favor a tectonic origin because of the lack of a seasonal pattern in seismicity. Instead, an oceanic origin is more favored because of the well-documented fact that oceanic wave activity and microseism excitations show similar seasonal pattern: more powerful during the local winter (e.g., Stehly et al., 2006; Stutzmann et al., 2009; Landès et al., 2010; Hillers et al., 2012; Reading et al., 2014). Next, we analyze the correlations between P_{dmc} signals and oceanographic data at a global scale.

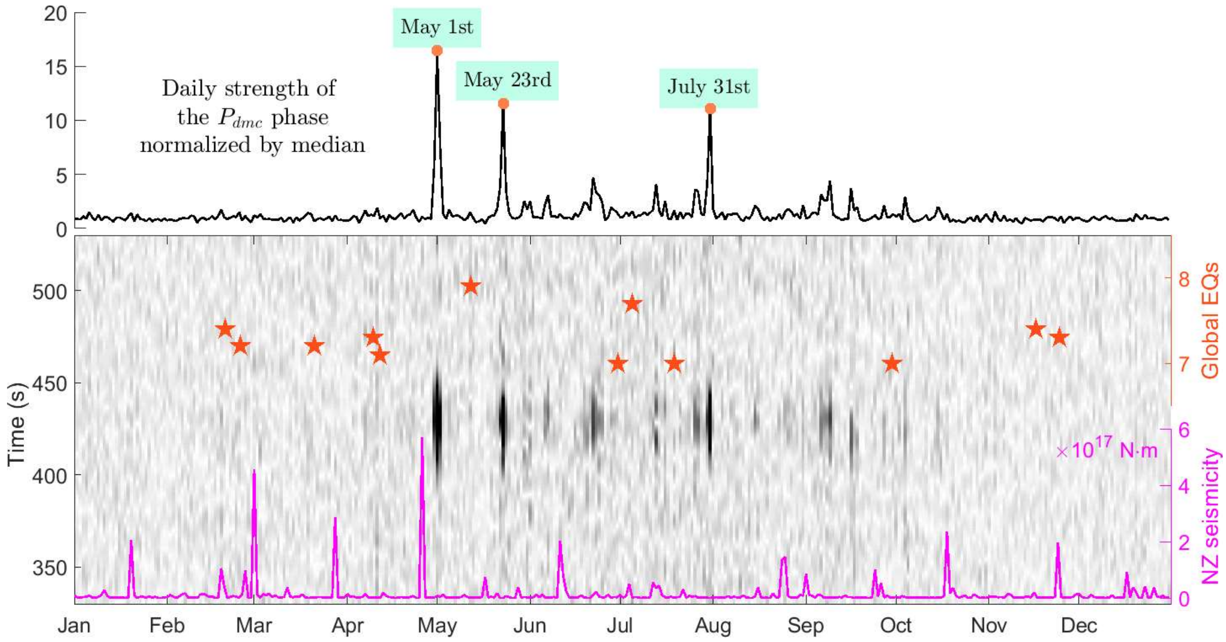

Figure 2. Temporal variations in the strength of daily P_{dmc} signals, in comparisons with the daily cumulative seismic moments in NZ (pink line at bottom; for earthquake magnitudes above 2.0 in GEONET catalogue) and global large earthquakes (stars; magnitudes above 7.0 in USGS catalogue; see Table S1 for a full list of earthquakes in 2008 above magnitude 5.5). The background image is composed of columns of daily envelopes of beamed FNET-LAPNET noise

correlations. Darker color represents larger amplitude. The curve on the top shows the daily P_{dmc} strength derived from the daily envelopes. Dates of the three largest peaks are labeled.

4 Correlation analysis

The sea state is composed of ocean waves at various frequencies and propagation directions. The nonlinear interaction between nearly equal-frequency ocean waves traveling in nearly opposite directions is equivalent to a vertical random pressure applied to the ocean surface (Longuet-Higgins, 1950; Hasselmann, 1963), so that microseisms are generated. Figure 3(a) shows a global map of average Power Spectral Density (PSD) of the equivalent surface pressure for a seismic period of 6.2 s, during the northern winter months of 2008. The hindcast PSD data are simulated by Ardhuin et al. (2011) and Rascle & Ardhuin (2013), based on the microseism excitation theory of Longuet-Higgins (1950) and Hasselmann (1963). The most energetic microseism excitations occur in the northern Atlantic south of Greenland and Iceland (near LAPNET), and in the northern Pacific between Japan and Alaska (near FNET). Figure 3(b) shows the map for the austral winter months, with the strongest excitations occurring between NZ and Antarctic (near GEONET). The seasonal pattern of oceanic microseism excitations results from the same pattern of global wave climate (Figs 3e-f). The seasonal pattern of the P_{dmc} strength agrees with that of the microseism excitation and wave climate in the effective source region south of NZ.

We compute the correlation coefficient (denoted as r) between the P_{dmc} strength and the source PSDs at each grid point, and thereby obtain a global correlation map (Fig. 3c). The largest r value for P_{dmc} and source PSD arises at [47°S, 177°E] in the effective source region (*E* in Fig. 3c). The corresponding time series of daily source PSDs is plotted in Fig. 4, in parallel with the P_{dmc} strength. Large peaks in the P_{dmc} series have good correspondence with large peaks in the source PSD series. From Fig. 3(c), one can observe a broad region of positive r values (red colors; roughly, south Atlantic, south Pacific, and Indian ocean). However, the positive correlation does not imply a causality between the P_{dmc} phase and the sources outside the effective region *E*. We ascribe the apparent positive correlation to the spatial correlation of the time-varying microseism excitation. As shown in Fig. 3(d), the source at [47°S, 177°E] in region *E* exhibits a similar pattern of apparent correlations with global sources as in Fig. 3(c). Despite the microseism excitations at varying locations are independent (Hasselmann, 1963), we note that the independence refers only to the phase information. The time variations of microseism source power are spatially associated. That is not surprising since the interacting ocean waves that excite microseisms could be driven by the same storms and swells can propagate freely over thousands of kilometers away (Ardhuin et al., 2009). We also notice there are high-r regions that may not be fully explained by the spatial association. These regions are characterized by low intensity of microseism excitations in Figs 3(a-b). A striking example is around [12°N, 88°E] in the Bay of Bengal (*F* in Fig. 3c). From Fig. 4, it can be seen that the source PSD series for [12°N, 88°E] is dominated by a single peak around May 1st, coincident with the largest P_{dmc} peak. This coincidence leads to a high value of correlation coefficient. However, the Bay of Bengal is far away from the FNET-LAPNET great circle, which is inconsistent with the source imaging shown later in Fig. 5. Thus, the high correlation is spurious and does not imply a causality relationship between the microseism sources in the Bay of Bengal and the P_{dmc} signals. Figures 3(g-h) show the correlation maps for h_s, which will be discussed later.

As shown in Fig. 4, prominent peaks in the P_{dmc} series have correspondence in the source PSD series for the effective source at [47°S, 177°E]. However, there are some peaks in the latter without correspondence in the former (see the labeled dates in Fig. 4b for examples). Note that here the P_{dmc} strength is compared to the microseism source PSD at single point in Fig. 4, whereas the effective sources spread over a region. One needs to verify if the peak disparities observed from Figs 4(a-b) can be ascribed to the neglect of the spreading of the effective source region. To evaluate an overall microseism excitation in the effective source region, the bathymetric effect on *P*-wave excitation should be considered (in the previous analysis for single point locations, the consideration of bathymetric effect is unnecessary because a scaling over the source PSD series does not change the value of the correlation coefficient between P_{dmc} and source PSD). Using the equations proposed by Gualtieri et al. (2014) and the bathymetry around NZ (Fig. 5a), we compute the bathymetric amplification factors for *P* waves at a period of 6.2 s (Fig. 5b; see Fig. S3 for comparisons between the factors calculated following Longuet-Higgins, 1950 and Gualtieri et al., 2014). The factors vary largely with locations. Also, note that the P_{dmc} phase has different sensitivity to the sources in the effective region, or say, the sources make varying contributions to the P_{dmc} signal. The power of sources should be weighted in the averaging. We obtain the weights by back-projecting the beam power of noise correlations onto a global grid (Fig. 5c; see Supplementary for technical details). Figure 5(d) shows the map of annually averaged source PSDs surrounding NZ and Fig. 5(e) shows the map after the modulation of the bathymetric amplification factors in Fig. 5(b). The spatial patterns are altered significantly, indicating the importance to account for the bathymetric effect. The final source imaging that has been weighted by Fig. 5(c), is plotted in Fig. 5(f). It agrees well with the effective source region *E* determined from the correlation map in Fig. 3(c). Replacing the annual PSD map in Fig. 5(d) with daily PSD maps, we obtain maps like Fig. 5(f) for each date. Averaging over the map leads to the time series of daily intensity in the effective source region (labeled as effective source intensity in Fig. 6). Averaging over a wide region has the advantage that the effects of potential source location errors due to the simplification of Earth model for fast travel time calculation, which have been addressed in some single array back-projection studies (e.g., Gal et al., 2015; Nishida & Takagi, 2016), can be largely reduced. From Fig. 6, one can see that the new effective source intensity series has almost the same peaks as the source PSD series for [47°S, 177°E] in Fig. 4(b), suggesting that the observed peak disparities are caused by other reasons. Next, we investigate if the disparities are caused by errors in the simulation of hindcast data or if there are other physical explanations.

The microseism source PSD data are simulated from the hindcast data of ocean wave directional spectra base on the excitation theory of Longuet-Higgins (1950) and Hasselmann (1963), which have no constrains from seismological observations. One should consider the accuracy of the simulation: can we ascribe the peak disparities in Fig. 4 to the simulation error or not? The seismic noise records from the GEONET array adjacent to the effective source region provide the opportunity to validate the simulation. To obtain the daily microseism noise levels at GEONET, we apply the Hampel filter, a variant of the classic median filter, to the continuous seismograms to discard earthquakes and anomalous impulses. The filter replaces outliers with the medians of the outliers' neighbors and retains the normal samples. Technical details are provided in section S4 of the Supplementary. The resultant GEONET noise level exhibits a good correlation with the effective source intensity (r = 0.7). We thus deem that the numerical simulations are statistically reliable. When the effective source intensity is high, the GEONET noise level should also be high (see the peaks marked by dots in Fig. 6 for examples). However, due to the great spatiotemporal variability of noise sources in the effective region and the complexity of seismic waves propagating from ocean to land (Ying et al., 2014; Gualtieri et al., 2015), a larger peak in the source intensity series does not necessarily imply a larger peak in the noise level time series (e.g., see diamonds in Fig. 6 for examples). We also emphasize that a high GEONET noise level does not need to always have a correspondence in the source intensity (see squares in Fig. 6 for example), because the GEONET stations record microseisms emanating from noise sources all around, not only from the effective source region.

The above analysis explains the observed disparities between the P_{dmc} strength and the effective source intensity. From Fig.6, one can see that the disparities primarily emerge in the shaded period when dominant microseism sources shift to the north hemisphere. The shading roughly separates the northern winter from the austral winter. The correlation between P_{dmc} strength and effective source intensity is low in the shaded period (r = 0.16), in contrast to the high correlation during the unshaded period (r = 0.74). Large P_{dmc} peaks always emerge on dates during the austral winter when the effective source intensity is much higher than its median, and meanwhile, noise levels at FNET and LAPNET are below their respective medians (see dots in Fig. 6 for examples). The seasonal variations of oceanic sources in the southern hemisphere are less strong than in the northern hemisphere (Fig. 3). On some dates (see triangles in Fig. 6 for examples), the effective source intensity can be considerable, but relevant P_{dmc} peaks are still missing. We notice that the corresponding microseism levels at FNET and LAPNET are obviously above their medians. Intensive ocean activity and microseism excitations in the north Pacific and Atlantic, lead to increased microseism noise levels at FNET and LAPNET. The P_{dmc} strength is anti-correlated with microseism noise levels at FNET (r = -0.12) and LAPNET (r = -0.18). We hereby conjecture that the microseism energy from the distant effective source region is dwarfed by the energetic microseisms excited by oceanic sources closer to the correlated FNET and LAPNET arrays, and consequently, P_{dmc} signals are overwhelmed by the background noise in the FNET-LAPNET cross-correlations. Last, we mention that the median threshold in Fig. 6 separates the major features of the time series described above, but there is no guarantee that it is a perfect threshold due to the nonlinear relationships between the P_{dmc} strength and the noise levels at the arrays.

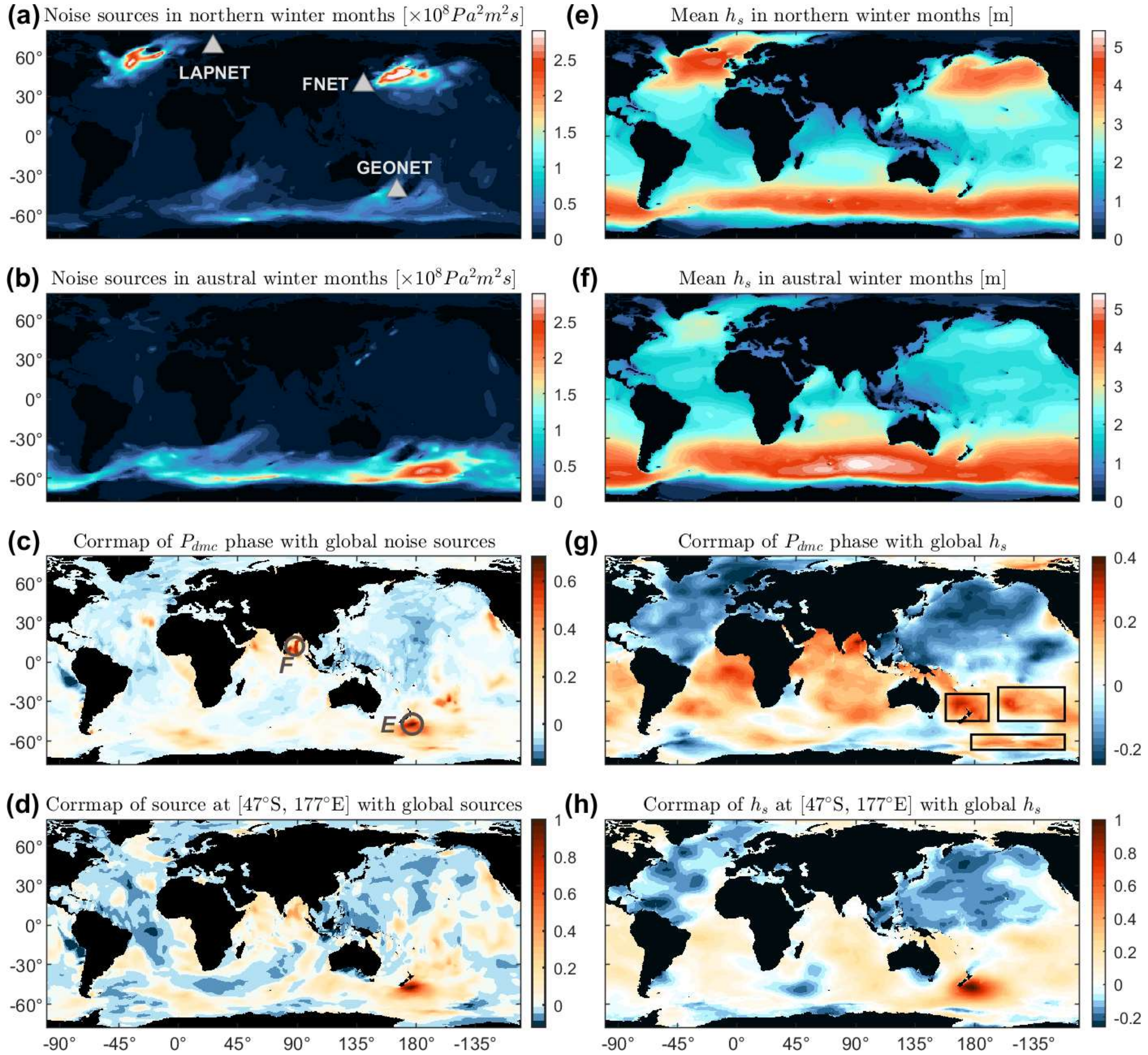

Figure 3. (a) Global map of average PSD of oceanic microseism sources in 2008 northern winter months (Jan. to Mar. and Oct. to Dec.), for a seismic perod of 6.2 s. (b) Similar to (a) but for 2008 austral winter months (Apr. to Sep.). (c) Correlation map (corrmap) for the P_{dmc} strength and global microseism noise sources. Circles mark two regions with highest correlation coefficients: *E*, effective source region surrounding [47°S, 177°E] south of NZ; *F*, fake highly-correlated region surrouding [12°N, 88°E] in the Bay of Bengal. (d) Correlation map for the source at [47°S, 177°E] and global sources. (e) Mean significant wave height (h_s; four times the square root of the zeroth-order moment of ocean-wave frequency spectrum) in northern winter months. (f) Similar to (e) but for austral winter months. (g) Correlation map for the P_{dmc} strength and global wave heights. (h) Correlation map for wave heights at [47°S, 177°E] and global wave heights. The oceanographical hindcast data are provided by the IOWAGA products (Rascle & Ardhuin, 2013).

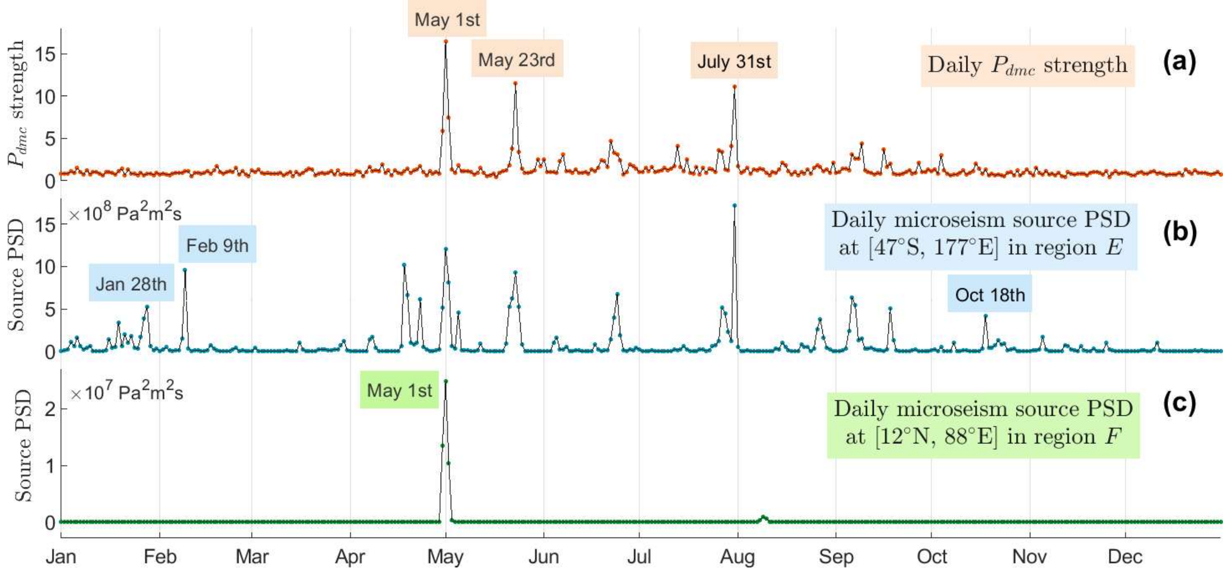

Figure 4. True correlation (r = 0.73) between (a) the P_{dmc} strength from Fig. 2 and (b) the power of source at [47°S, 177°E] in the effective source region (*E* in Fig. 3c), and spurious correlation (r = 0.71) between P_{dmc} and (c) the power of source at [12°N, 88°E] in the Bay of Bengal (*F* in Fig. 3c).

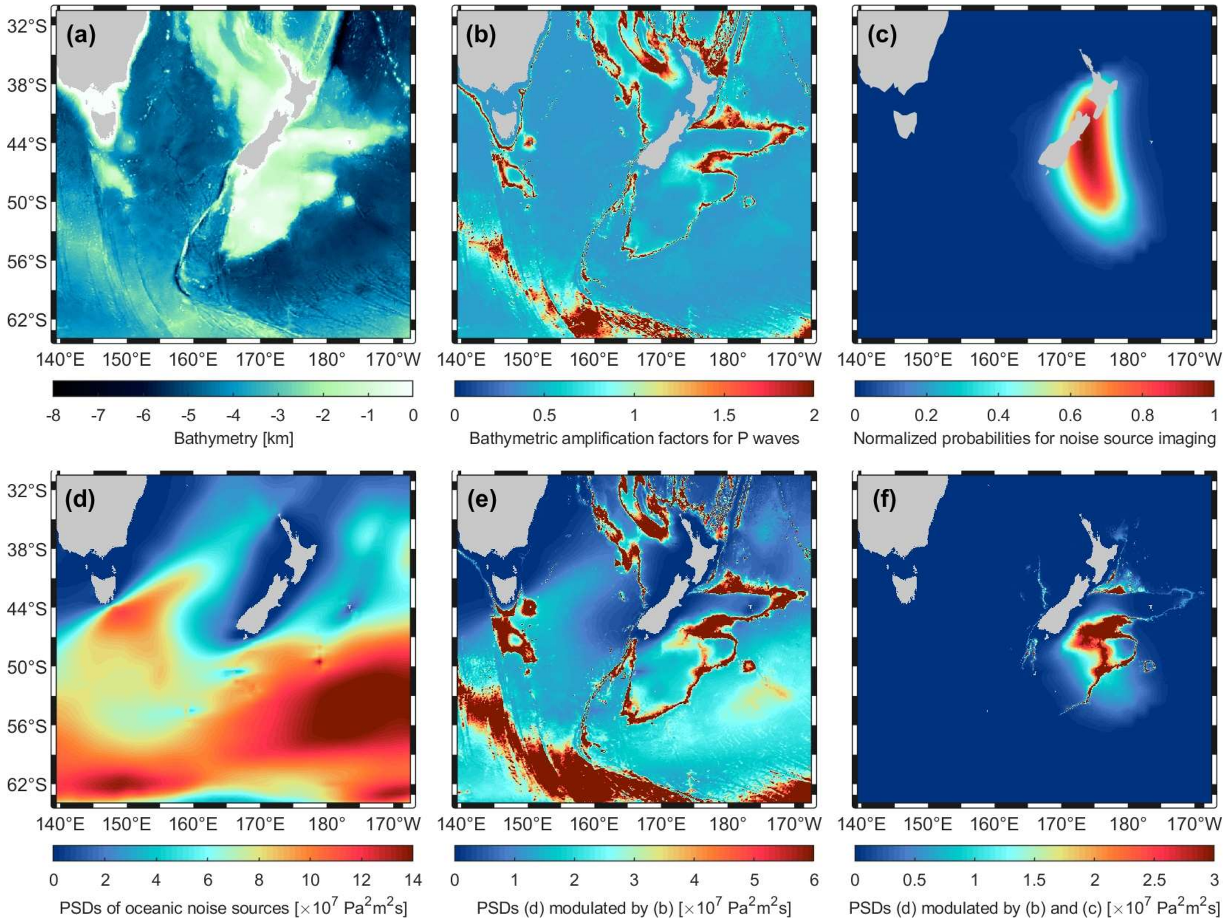

Figure 5. (a) Bathymetry around NZ. (b) Bathymetric amplification factors for *P*-type waves. (c) Imaging of effective sources obtained from the back-projection of the FNET-LAPNET noise correlations. (d) Annual average of source PSDs in 2008. (e) Source PSDs in (d) modulated by the factors in (b). (f) Source PSDs in (e) further modulated by the weights in (c).

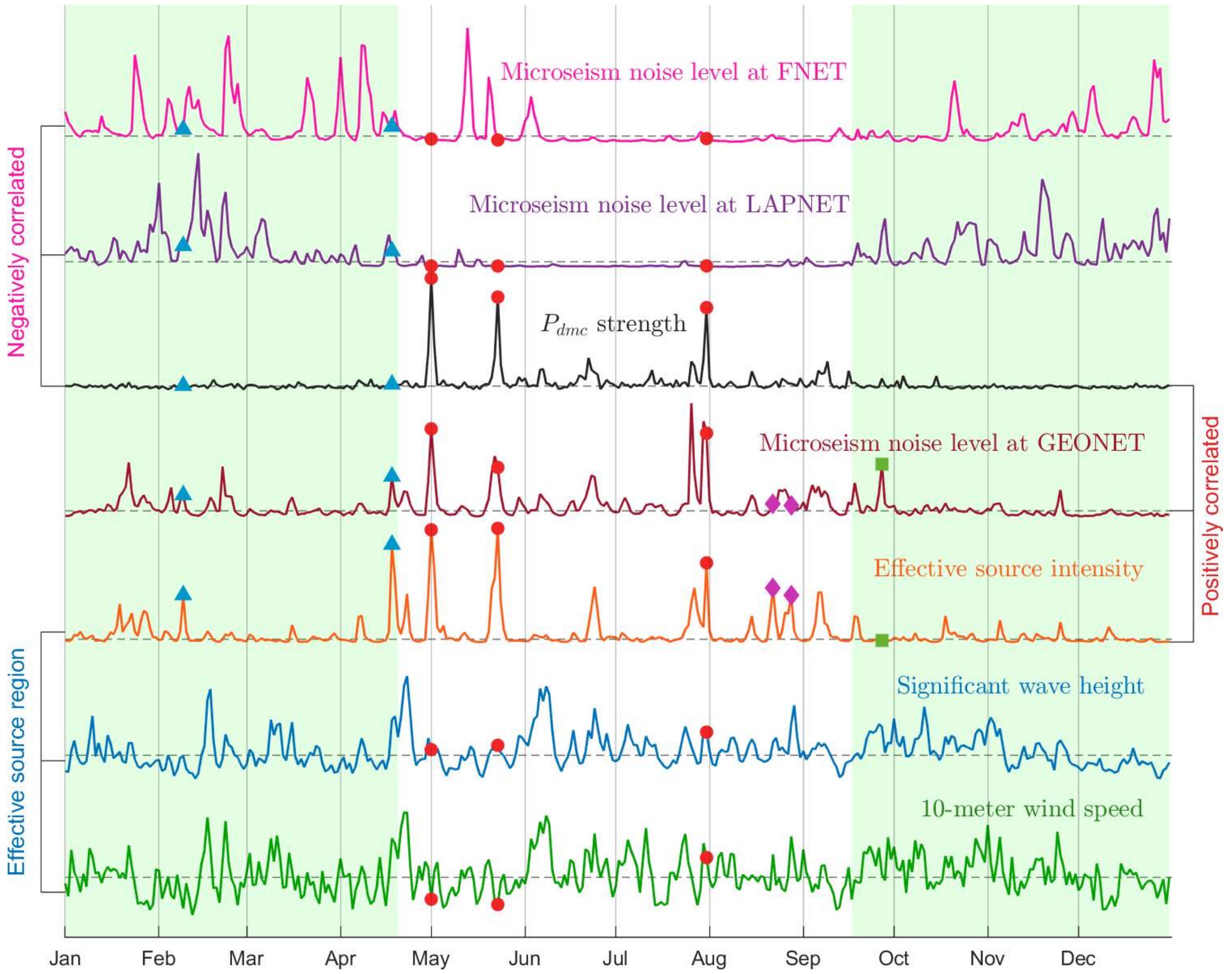

Figure 6. Temporal variations of daily P_{dmc} strength, microseism noise levels at three networks, and average wind speeds, wave heights and microseism excitations in the effective source region. The curves are normalized by their own maximums. Dashed horizontal lines denote their respective medians. Symbols mark some dates cited in the main text. When computing the effective source intensity, the bathymetric factors in Fig. 5(b) and weights in Fig. 5(c) are used. When computing the average wind speeds and wave heights, weights in Fig. 5(c) are used.

5 Discussions and conclusions

In this study, we explore the relations between the noise-derived P_{dmc} signals and global oceanic microseism sources using spatiotemporal correlation analysis. The effective source region E for the P_{dmc} phase is successfully identified from the correlation map in Fig. 3(c), which is consistent with that determined from the seismological back-projection in Fig. 5(c). The correlation map provides a convenient way to identify the effective sources of noise-derived seismic signals.

In our case, the seismic networks used for noise correlation are located in the northern hemisphere, while the effective source region is in the southern hemisphere. Ideally, we expect a correlation map with the following features: positive correlation with sources in the effective region, and negative or insignificant correlations with other inefficient sources. Positive correlation indicates a contribution to the construction of P_{dmc} signal from noise correlations, negative correlation implies an adverse impact, and insignificant correlation (decorrelation) means a negligible effect on the signal construction. However, we obtained a correlation map roughly showing that, the P_{dmc} signal is correlated with the southern sources and anti-correlated with the northern sources. The correlation with southern sources outside the effective region can be interpreted with the spatiotemporal correlation of the power of the microseism sources in the southern oceans, due to the large span of ocean storms and the long-range propagation of swells. The anti-correlation with the northern sources, can partly be explained by the well-known reverse seasonal patterns of oceanic microseism excitations in the south and north hemispheres (Stutzmann et al., 2009; Landès et al., 2010; Hillers et al., 2012; Reading et al., 2014). Another important reason is that compared to the remote effective sources in the south hemisphere, the northern sources closer to the correlated stations have larger impacts on the microseism noise levels at stations. Strong energy flux from the northern sources outshines the microseism energy coming from the distant effective sources. That deteriorates the construction of the P_{dmc} phase. The noise-derived P_{dmc} signals are primarily observable in the austral winter. That can be, on one hand, attributed to the stronger effective source intensity during that period, and on the other hand, to the relative tranquility in the northern oceans.

In Fig. 7, we summarize the classification of noise sources, the decomposition of wavefields, and the associations to the constituents of the inter-receiver noise correlation function. The diagram of Fig. 7(a) explains the relationships using the case study of the P_{dmc} phase discussed above. We generalize Fig. 7(a) to the derivation of an arbitrary signal (referred to as the target signal for convenience) from ambient noise wavefields (Fig. 7b). The noise correlation function is composed of the target signal, any other signals and background noise. A source or a wave is called effective if it contributes to the construction of the target signal from noise correlations. Otherwise, it is called ineffective. The construction of the target signal is exclusively ascribed to the interference between the effective waves. Stronger effective sources (relative to ineffective sources) imply more effective waves in the total wavefield, and thereby, a better quality for the noise-derived target signal. Note that not all waves emanating from the effective sources, but only those following specific ray paths, are effective. There might be multiple pairs of seismic phases that could contribute to the construction of the target signal. However, their relative strength matters. As for the case of the P_{dmc} phase, the effective waves are *P* and *PKPab*, which are both prominent phases in the ballistic wavefield. Li et al. (2020) showed that the *PcP-PKPab* correlation and the *PcS-PcPPcP* correlation, could also lead to a signal at around the P_{dmc} emerging time. However, the *PcP*, *PcS*, and *PcPPcP* waves are weak phases in the ballistic wavefield, and thereby have minor contributions to the P_{dmc} signals. We emphasize that the sketch in Fig. 7(b) is only suitable for the ambient noise wavefields that are dominated by ballistic waves.

From Fig. 6, one can observe a high correlation between wind speed and wave height in region *E* ($r = 0.74$). It indicates that the ocean waves in region *E* are likely dominated by the waves forced by local winds. The correlation between wave height and microseism excitation is low ($r = 0.25$), implying a dominant role of the freely propagating swells in exciting the microseisms. Extreme sea state does not guarantee strong microseism excitation. That is not surprising according to the microseism excitation theory (Hasselmann, 1963; Longuet-Higgins, 1950): the excitation is proportional to the product of the heights of the colliding equal-frequency ocean waves. In lack of equal-frequency waves coming from opposite directions, even extreme wave climate cannot incite strong secondary microseisms. In contrast, for large peaks in the microseism excitation, the corresponding wave heights are generally moderate (e.g., on May 1st and 23rd). On these two dates, the low wind speeds but moderate wave heights in region *E* suggest that the ocean waves are dominantly the freely travelling swells from elsewhere, as also illustrated in the supplementary movie S1. Oppositely propagating equal-frequency swells collide with each other and incite strong microseisms. Our analysis and observations agree with those of Obrebski et al. (2012), who investigated a specific case that small swells from two storms meeting in the eastern Pacific generate loud microseism noise. There are also examples showing that wind waves can play a role in the excitation of microseisms, for instance, around July 31st when the local winds, wave height, and microseism excitations are all strong. Such examples are few. The good consistency between the temporal variations in the P_{dmc} strength, the effective source intensity and the NZ microseism noise level (Fig. 6), provides extra supports to the analysis of the P_{dmc} observations and the quasi-stationary phase arguments proposed by Li et al. (2020). It also gives credits to the validity of the numerical modeling of oceanic microseism sources by Ardhuin et al. (2011) and Rascle & Ardhuin (2013).

We have described above the implications of this study in seismology and in understanding the process of microseism excitation. Now, we discuss the significance in oceanography and climate science. Well-documented historical ocean storms and wave climate are valuable for improving our understanding of climate change and global warming (Ebeling 2012). However, modern satellite observations of ocean waves and storms have a history of merely decades. Microseisms are induced by storm-driven ocean waves (Ardhuin et al., 2015; Hasselmann, 1963; Longuet-Higgins, 1950). The records of microseisms contain the imprint of climate (Aster et al., 2010; Stutzmann et al., 2009). Instrumental observation of microseisms has an over-century history, and started much earlier than the modern observations of ocean waves and storms. It has been a long-lasting effort for the seismological community to digitalize the historical analog seismograms (Bogiatzis & Ishii, 2016; Lecocq et al., 2020). Researchers expect that past seismic records can be used to recover undocumented historical ocean storms and wave climate (Ebeling 2012; Lecocq et al., 2020).

This study confirms that it is possible to detect remote microseism events (burst of microseism energy) with land observation of microseisms. We demonstrate that the noise-derived P_{dmc} signals can be employed to monitor microseism events in a specific ocean region (Fig. 5). The remote monitoring of microseisms is promising as an aid to improving wave hindcast, in similar manners as demonstrated by Stopa et al. (2019). The comparative analysis in Fig. 6 indicates that the remote event detection could be effective in the absence of strong sources near the stations, otherwise the detection could fail. Stations at low latitudes where wave climate and microseism excitation are relatively mild, or inland stations far from oceans, should have better performance in remote monitoring.

Energetic microseism excitation does not always need extreme in situ wave heights, and extreme wave heights do not necessarily produce powerful microseisms (Obrebski et al., 2012; and this study). It imply that secondary microseism events are not a perfect proxy for the extremal in situ wave climate. However, it does not mean the long-lasting attempt to monitor remote sea state and ocean storms with land observation of secondary microseisms is futile. In the P_{dmc}-h_s correlation map (Fig. 3g), the largest *r* values do not fall in the effective region *E* as in the P_{dmc}-source correlation map (Fig. 3c), but in surrounding regions with moderate to high ocean wave activity (the bounded areas in Fig. 3g). We speculate that these regions could be the birthplaces of the colliding swells that generate the secondary microseisms in region *E*, or the ocean waves in these regions are driven by the same storms as the colliding waves in region E (see the spatial links of h_s from Fig. 3h and supplementary movie S1). The detection of a microseism event could affirm the existence of the causative storms that generated the ocean waves propagating to the location of the microseism event, although the storms could be distant from the events.

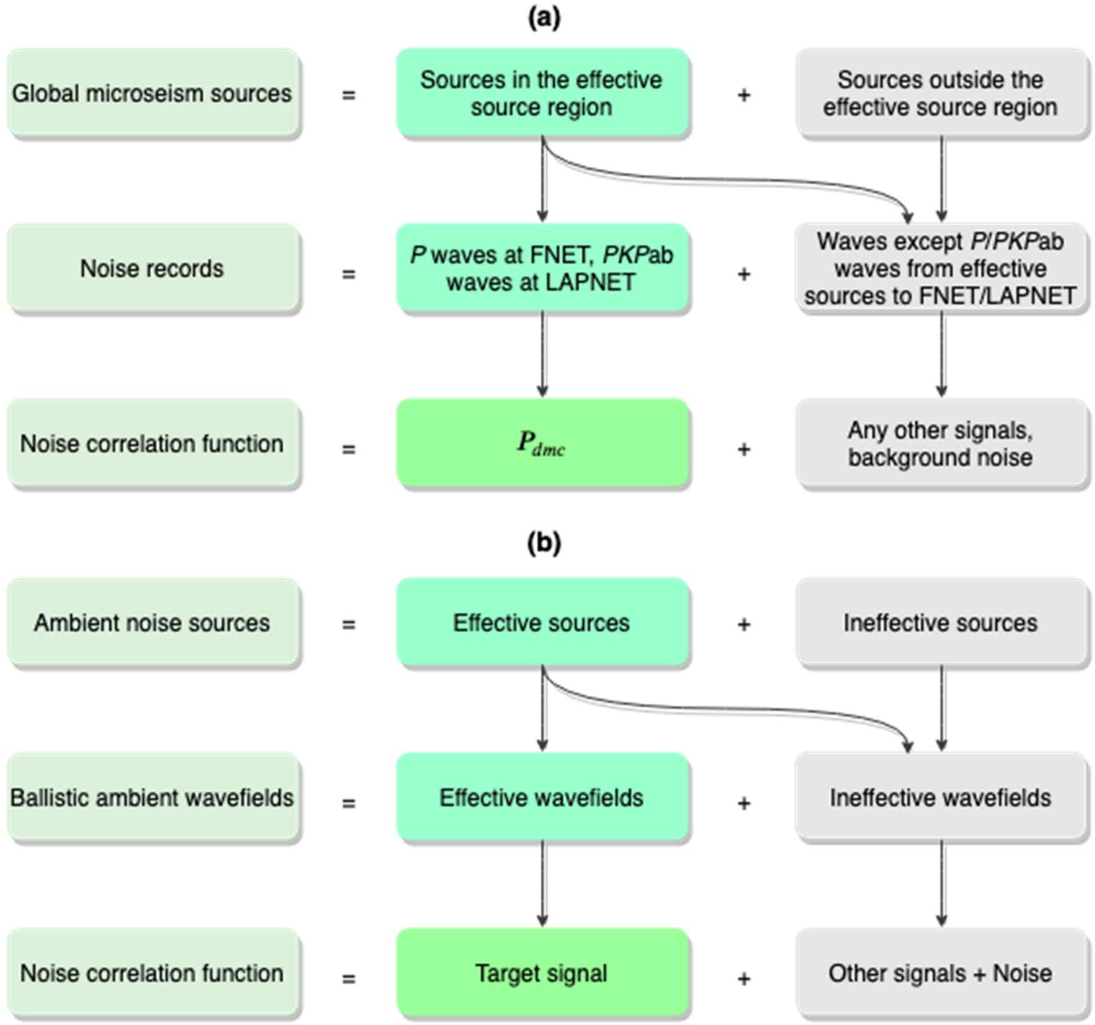

Figure 7. (a) Sketch explanation for the relationships between microseism noise sources and the noise-derived P_{dmc} signal. (b) Generalization of diagram (a) for an arbitrary signal derived from ambient noise wavefields that are dominated by ballistic waves.

Acknowledgments and data

The seismic data of FNET and LAPNET were provided by the National Research Institute for Earth Science and Disaster Resilience (http://www.fnet.bosai.go.jp/; last access: June 2018) and the Réseau Sismologique & Géodésique Français (http://www.resif.fr/; last access: June 2018), respectively. The seismic data of GEONET and the earthquake catalogue of New Zealand were provided by the GEONET Data Center (https://www.geonet.org.nz/; last access: June 2018). The global earthquake catalogue was provided by the U.S. Geological Survey (https://earthquake.usgs.gov/; last access: June 2018). The wind hindcast data were provided by the European Centre for Medium-Range Weather Forecasts (https://www.ecmwf.int/; last access: June 2018). The hindcast data of wave heights and microseism source PSDs were provided by the IOWAGA products (Rascle & Ardhuin, 2013). The bathymetry data were extracted from ETOPO1 Global Relief Model (Amante & Eakins, 2009). The computations were performed mainly on the ISTerre cluster. This work was supported by Labex OSUG@2020 (Investissements d'avenir-ANR10LABX56) and the Simone and Cino del Duca Foundation, Insitut de France. The authors acknowledge the support from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (grant agreement No 742335, F-IMAGE). We also acknowledge Anya M. Reading and an anonymous reviewer for their valuable reviews that helped to improve the clarity of our paper, and special thanks to the editor for allowing the deadline extension during the COVID-19 pandemic.

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[Geochemistry, Geophysics, Geosystems]

Supporting Information for

Spatiotemporal correlation analysis of noise-derived seismic body waves with ocean wave climate and microseism sources

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Contents of this file

Additional Supporting Information (Files uploaded separately)

Caption for Table S1
Caption for Movie S1

Text S1. FNET-LAPNET noise correlations

The correlation function C_{AB} between two seismograms (S_A and S_B) is given by

$$C_{AB}(\tau) = \frac{\sum_i S_A(i) S_B(i-\tau)}{\sqrt{\sum_i S_A^2(i) \sum_i S_B^2(i)}}. \qquad \text{(S1)}$$

The resultant C_{AB} consists of an acausal part and a causal part, that correspond to the negative lags ($\tau < 0$) and the positive lags ($\tau > 0$), respectively. For efficiency, it is routine to compute the correlation function with the Fast Fourier Transform:

$$C_{AB}(\tau) = \frac{\mathcal{F}^{-1}[\mathcal{F}(S_A)\mathcal{F}^*(S_B)]}{\sqrt{\sum_i S_A^2(i) \sum_i S_B^2(i)}}. \qquad \text{(S2)}$$

Figure S1(a) shows the acausal and causal sections of FNET-LAPNET noise correlations in 2008 that are filtered between 5 s and 10 s and binned in distance intervals of 0.1°. The acausal section is flipped to share the time axis with the causal section. The expected locations of the acausal and causal noise sources are marked by stars on the maps of global microseism source PSDs and ocean wave heights in Fig. S1(b). The ocean wave activities and microseism excitations at the acausal source region are intense, while those in the causal source region are fainter. Consequently, the P_{dmc} phase is only observable from the acausal noise correlations.

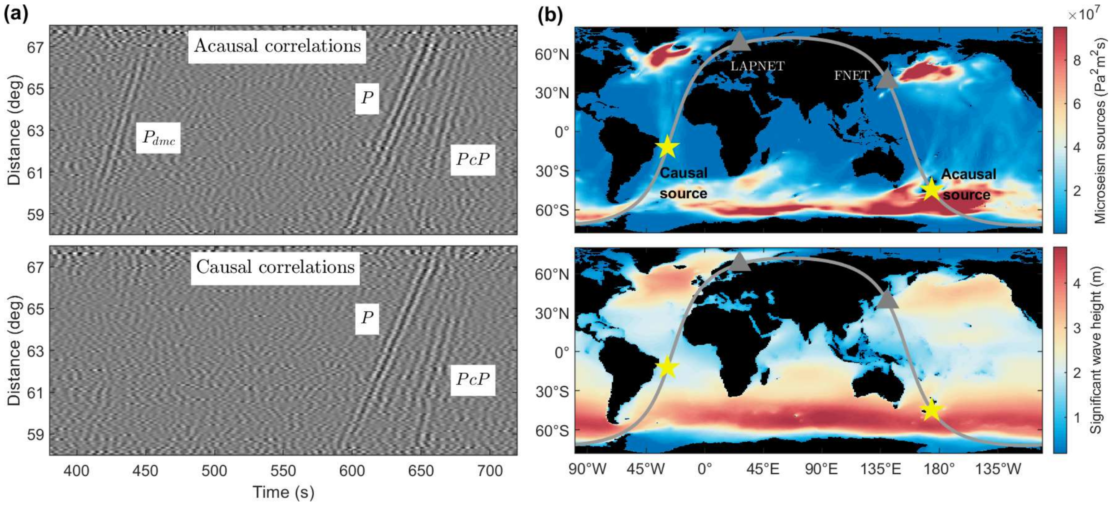

Figure S1. (a) Acausal and causal sections of FNET-LAPNET noise correlations in 2008. (b) Global maps of 6.2 s period secondary microseism sources and significant wave heights in 2008.

Text S2. Noise source imaging by back-projection

Assuming the interferometry between *P* waves at FNET and *PKPab* waves at LAPNET, we image the effective noise sources through the back-projection of the FNET-LAPNET noise correlations. We beam the FNET-LAPNET noise correlations and assign the beam power

$$P_s = \langle\langle\, C_{ij}\big(t + t_{si} - t_{sj}\big)\rangle_{ij}^2\rangle_t, \qquad \text{(S3)}$$

onto a 0.5° × 0.5° grid as the probabilities of noise sources on the global surface. In the above equation, $\langle\cdot\rangle_x$ means the average over x, C_{ij} is the correlation function between the i-th FNET station and the j-th LAPNET station, t_{si} is the traveltime of the *P* wave from the s-th grid point to the i-th FNET station, and t_{sj} is the traveltime of the *PKPab* waves from the s-th source to the j-th LAPNET station. The inter-station noise correlations are windowed before the beamforming (Fig. S2a). The noise source imaging for the annually stacked noise correlations is plotted in Fig. S2(c). Only the region surrounding NZ is shown. Outside the region, hardly can the *P* wave reach FNET or the *PKPab* waves reach LAPNET. Besides a well-focused imaging of the expected source region in the ocean south of NZ, we notice a secondary spot to the west. In comparisons with the power map of oceanic microseism noise sources in Fig. 5(e), we ascribe it to the strong microseism excitation in the ocean south of Tasmania. We also back-project the daily noise correlations on 2008-05-01 (Fig. S2b), when the P_{dmc} phase reaches the largest strength through the year (Fig. 2). As shown in Fig. S2(d), an exclusive source region is imaged, which agrees with the dominant spot in Fig. S2(c).

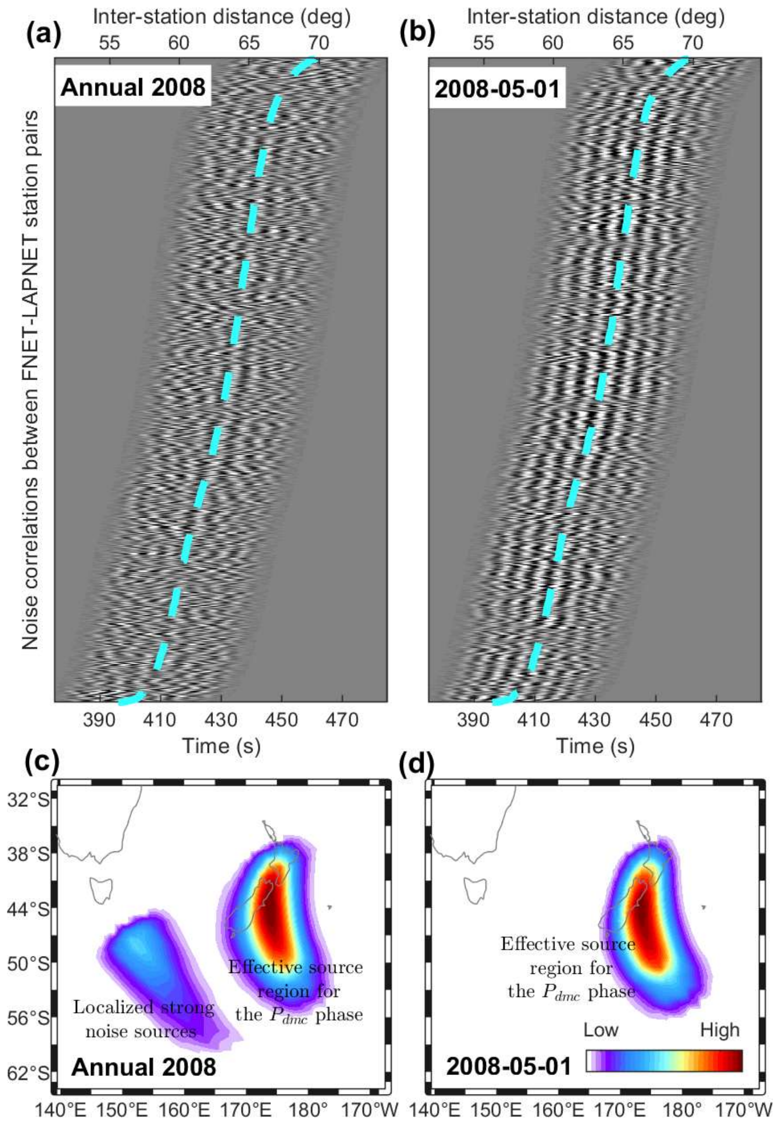

Figure S2. Inter-receiver noise correlations for all FNET-LAPNET station pairs: (a) stacked over the year of 2008; (b) on single day of 2008-05-01. The waveforms are windowed around the P_{dmc} phase. Dashed lines indicate inter-station distances. Back-projection imaging of noise sources: (c) using data from (a); (d) using data from (b).

Text S3. Bathymetric amplification factors

Figure S3 compares the bathymetric amplification factors surrounding New Zealand for *P* waves and Rayleigh waves. The factors for *P* waves are computed using the equations proposed by Gualtieri et al. (2014), for a seismic period of 6.2 s and a slowness of 4.6 s/deg. The factors for 6.2 s period Rayleigh waves are obtained by interpolating the table given by Longuet-Higgins (1950).

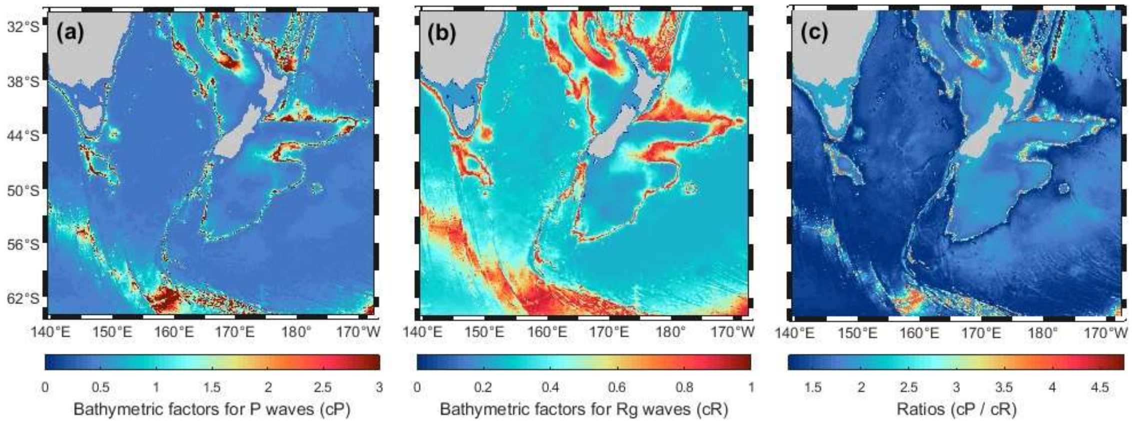

Figure S3. Bathymetric amplification factors for (a) *P* waves and (b) Rayleigh waves. (c) Ratios between the factors for *P* waves and for Rayleigh waves.

Text S4. Microseism noise levels at seismic networks

The continuous seismograms record not only the background vibrations of Earth, but also ground motions induced by seismicity or other events. Instrumental malfunction also leads to anomalous (e.g., nearly vanishing or extremely large) amplitudes in the records. These extreme amplitudes (outliers) could bias the estimates of microseism noise power. It is necessary to get rid of them from the ambient noise records before the computation of noise power. Mean and median filters are the common tools for this task. However, they modify all the samples. Here, we prefer to use a variant of the median filter called Hampel filter. In contrast to the median filter that replace all samples with local medians, the Hampel filter detects outliers by compare a sample with the neighboring samples. A sample is replaced by the local median if it deviates k times of the median absolute deviation (MAD) from the local median, or else, it is unchanged.

We filter the vertical components of the continuous seismograms around 6.2 s period. The seismograms are then divided into 15-min segments and the power of segments is computed. We apply the Hampel filter to the time series of noise power recursively. For each sample, we compute the local median and MAD of its eight neighbors (four before and four after). A sample is replaced by the median if it deviates from the median over three times of the MAD. The de-spiked time series is resampled from a 15-min interval to a 1-hour interval, by averaging over every four samples. Then, we apply the Hampel filter again and resample the time series to a 24-hour interval. The averaging of noise levels over all stations of a seismic network leads to the time series of array noise level. Before the averaging, the Hampel filter is applied again, to discard possible anomalous values at some stations (see Fig. S3 for the example of GEONET). The final time series of microseism noise levels for networks FNET, LAPNET and GEONET are shown in Fig. 6.

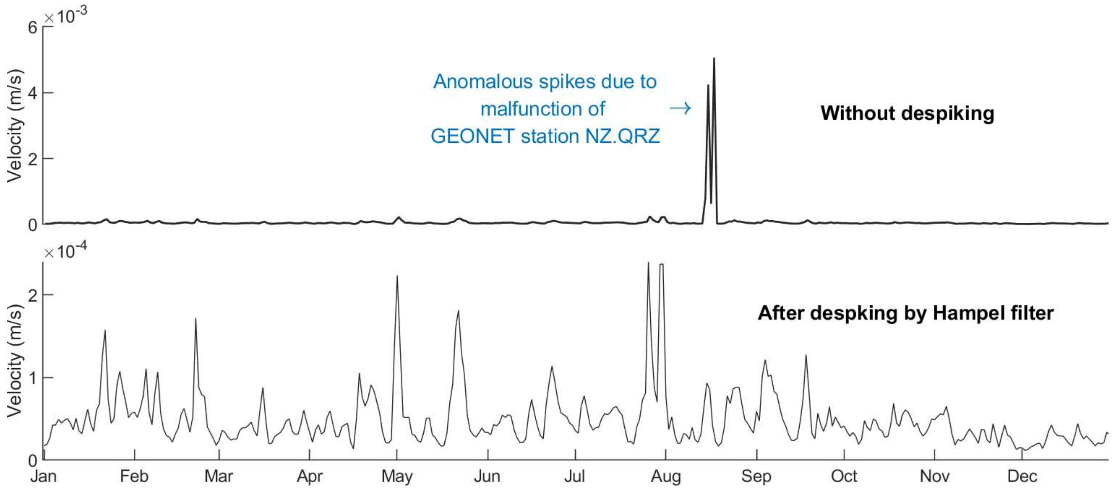

Figure S4. Comparison between the time series of daily GEONET noise levels with (lower) and without (upper) despiking using the Hampel filter.

Table S1. List of earthquakes (magnitude above 5.5) in 2008 extracted from the USGS catalogue, as a supplementary to the comparison between seismicity and P_{dmc} in Fig. 2 of the main text**.** On some dates with earthquakes near the FNET-LAPNET great circle (e.g., events 2008-08-25T11:25:19.310 and 2008-11-21T07:05:34.940), no large P_{dmc} is present, indicating that P_{dmc} is unrelated to earthquakes.

Movie S1. Daily evolutions of winds, ocean wave heights, and secondary microseism source PSDs around New Zealand in 2008. The closed lines superposing the upper panels depict the contour values of 0.1, 0.5, and 0.9 for the weights shown in Fig. 5(c). The source PSDs are modulated by the bathymetric factors shown in Fig. 5(e). In the bottom panel, the time series for the P_{dmc} strength and the weighted averages of the source PSD, wave height, and wind speed in the effective source region, are the same as those in Fig. 6 in the main text. See captions of Figs 5 and 6 for more details.